# Emergent Properties of Tumor Microenvironment in a Real-Life Model of Multicell Tumor Spheroids

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#### Abstract

Multicellular tumor spheroids are an important in vitro model of the pre-vascular phase of solid tumors, for sizes well below the diagnostic limit: therefore a biophysical model of spheroids has the ability to shed light on the internal workings and organization of tumors at a critical phase of their development. To this end, we have developed a computer program that integrates the behavior of individual cells and their interactions with other cells and the surrounding environment. It is based on a quantitative description of metabolism, growth, proliferation and death of single tumor cells, and on equations that model biochemical and mechanical cell-cell and cell-environment interactions. The program reproduces existing experimental data on spheroids, and yields unique views of their microenvironment. Simulations show complex internal flows and motions of nutrients, metabolites and cells, that are otherwise unobservable with current experimental techniques, and give novel clues on tumor development and strong hints for future therapies.

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#### Introduction

Multicellular tumor spheroids (MTS) stand out as the most important in vitro model of pre-vascular solid tumors [1-8]. MTS often have a regular, almost spherical structure, and their apparent simplicity has led to repeated attempts to capture their features with neat mathematical models. However, the absence of vascularization and the near sphericity hide an internal complexity which is not easy to tame either with analytic mathematical models [9-12], or with numerical models based on rough simplifications of the biological settings such as cellular automata or other latticebased models [13-16]. Moreover the presence of a growing necrotic core [1] and of an extracellular matrix [17], the appearance of convective cell motions [18], and the heterogeneous response to chemotherapics [19], point to the importance of MTS as an in vitro model of tumors, and most of all to their relevance to understand tumor heterogeneity, but they also point to the difficulties of producing a useful, predictive model of MTS.

The appearance of widely different resistance phenomena to antitumor therapies in similarly grown, isolated MTS of the same cell type [19] indicates that random fluctuation phenomena play an all-important role in the growth kinetics of MTS. It is wellknown that the discrete events at the single-cell level (like transitions from one cell-cycle phase to the next, mitosis, cell death, etc.) do display some randomness, and one can pinpoint the source of large-scale variability on these fluctuations, as they are amplified and propagated by cell-cell and cell-environment interactions. Thus, the complexity of MTS development can only be captured by a fine-grained, multiscale model, and we need a mathematical description at the single-cell level. Since cells communicate with other cells and the environment, the other actors of this complex play are the concentration gradients of important molecular species that depend on the structure of the extracellular space and of the facilitated transport processes into and out of individual cells, and the mechanical forces that push and pull cells as they proliferate with repeated mitoses and then shrink after death [20]. These processes mix with complex nonlinear interactions between the biochemical and the mechanical part, and this highlights again the importance of an effective model at the single-cell level.

On the basis of such motivations, we have developed a numerical model of MTS that incorporates a working model of single cells [21,22]. We have first put forward a broad outline of its structure in reference [23], and it differs from other models developed in the past [9–16] because it captures at the same time both the basic features of cell metabolism, growth, proliferation and death, and provides a true lattice-free calculation of cell motions, as they are pushed and pulled by the forces exerted by dividing cells, by the growth of other cells, and by the shrinking of dead cells. We also wish to stress that the model parameters are either derived from experiment or are deduced from reasonable theoretical arguments, so that, essentially, there are no free parameters – there can only be some residual variability in biophysically meaningful ranges – the model is truly predictive, and the results are not merely qualitative but quantitative as well.

Here we illustrate in broad terms the structure of the program and report the results of the first simulations of single spheroids (technical implementation details are relegated to Text S1). We find that the simulations agree quite well with experimental measurements on real spheroids, and show unexpected and important internal patterns. Moreover, we wish to stress that the methods delineated in this paper represent very general practical solutions to problems that are common to any simulation of cell clusters, and they are just as important.

#### Biochemical behavior of individual cells

The elementary building blocks in this model of MTS are the individual tumor cells that behave as partly stochastic automata [21,22]. Figure 1 summarizes the biochemical pathways that are included in the single-cell model: cell metabolism is driven by oxygen, glucose and glutamine, and transforms these substances into energy molecules, molecular building blocks and waste products, following the well-known biochemical reaction chains [24]. Further details can be found in the original papers [21,22] and in Text S1, which also includes important upgrades to the original model [21,22].

In the present version of the program, the stochasticity is mostly concentrated in the discrete events: for instance, mitochondria are partitioned at random between daughter cells at mitosis, and cells can die because of metabolite accretion, according to a Poissonian cytotoxicity model (see Text S1).

We remark that in this approach glutamine also stands for the wider class of aminoacids, and lactate is the paradigm of all metabolites: we use the concentrations of glutamine and lactate to represent these two classes of substances in phenomenological parameterizations wherever needed. Similarly we use the number of mitochondria and ATP content to model the dynamics of cell volume; the single-cell model also contains representative members of the cyclin protein class to compute the passage of checkpoints and entry into the different cell phases [21,22,25,26], and finally into mitosis (see also figure S1 for a sketch of the cell cycle in the simulation program).

The complete map of the biochemical pathways included in the simulation program is shown in figure S2. This map comprises only the most basic pathways, however we cannot afford to introduce a more complex network at this stage of program development. Indeed, our final aim is the simulation of MTS with a volume as large as 1 mm<sup>3</sup>, which corresponds to more than one million cells, so that simulation results overlap actual experimental measurements [19,27,28]. Since the differential system involves 19 independent biochemical variables per cell, we must eventually integrate at least 19 million coupled nonlinear differential equations for the biochemical cell variables alone (this grows to at least 25 million equations when we include the position and velocity variables), and thus even this minimal single-cell model leads to a daunting computational task (see Text S1 for further details on the algorithmic complexity of the program).

#### Reaction-diffusion processes and the environment

Substances like oxygen are transported into and out of cells by normal diffusion while molecules like glucose require facilitated diffusion processes. This means that cell membranes play an important role for substances like glucose, and that in this case the diffusion of each such molecular species towards cells in the inside of a spheroid needs the free volume in the extracellular space to proceed, and that we must model this space as well as the cells to obtain a realistic simulation. We have shown how to do this in reference [29], where we have also discussed ways to tame the



Figure 1. Rough sketch of the biochemical pathways incorporated in the model of single cells. We take into account the main metabolic pathways (glycolysis, oxidative phosphorylation through the TCA cycle and gluconeogenesis), including the role of mitochondria in the production of ATP. The model also includes protein and DNA synthesis, and checkpoints controlled by representative members of the cyclin family. The single-cell model has two spatial compartments (the inside of the cell and its immediate neighborhood, the extracellular space that surrounds it) and transport of substances between these compartments is regulated by transporters on the cell membrane that are also included in the model. The extracellular space of each cell communicates by simple diffusion with the neighboring extracellular spaces and with the environment. The complete map of the biochemical pathways is shown in figure S2. doi:10.1371/journal.pone.0013942.g001

exceedingly high stiffness of the very large set of reaction-diffusion and transport equations that arise in this context (see also Text S1). The external environment itself is included in these equations, and evolves in time as well. In the present model, each cell contributes 15 internal variables and 4 extracellular variables: these extracellular variables are the masses of oxygen, glucose, glutamine and lactate in the extracellular volume surrounding the cell. Because of its smallness, the extracellular space has an extremely short characteristic filling time, which can be as fast as few tens of microseconds. On the other hand, the macroscopic features of MTS evolve over times as long as months (i.e., times of the order of  $10^7$ s), and thus the numerical integrator must be able to handle phenomena that span 12 orders of magnitude in time [29]. The internal biochemical reactions included in the numerical model are much slower and their fastest characteristic times are only as low as 0:1s, much longer than the diffusion times [29,30]. The topology of diffusion in the extracellular spaces is obviously dictated by the cells themselves, and the program uses the network of cells centers as the scaffolding for the corresponding discretized diffusion problem. The links between the cells' centers - i.e., the proximity relations - are provided by a Delaunay triangulation [31,32], which is computed repeatedly [33] as the cluster of cells grows and rearranges itself under the pushes and pulls of volume growth, mitosis, and the shrinking of dead cells (see also figure S3). Moreover, the proliferation of cells means that both the number of cells and the total number of links steadily grow, and that the differential system of equations that model metabolism, transport and diffusion changes all the time, and becomes increasingly complex. The 3D Delaunay triangulation itself is not an exceedingly heavy computational burden for the program, as it turns out that efficient algorithms can compute it, on average, with O(N) time computational complexity [33–35], so that this algorithm is indeed feasible for very large clusters of cells.

#### Biomechanical evolution of the simulated MTS

Real cells have passive viscoelastic mechanical features, but they also move actively under the pushes of their own cytoskeleton, and to the best of our knowledge there is no comprehensive model of cellular biomechanics [36,37]. Thus, we resort once again to phenomenological simplifications, and the first and foremost is that our cells are stretchable spheres, characterized by their radius, and a few other parameters that specify their viscoelastic properties (see Text S1 for a more detailed description and the list of parameters). We also specify a pairwise interaction force between cells, repulsive when a cell pushes against a neighbor, and attractive when we try to detach it from its neighbor. For small deviations from the equilibrium distance, we assume that the interaction force is described by the Hertz model (explained in Text S1), while for large deformations due to compression we set the force to a fixed saturation value, and for large distances the attractive force decays to zero (see figure S4). The description of the interaction forces is tuned to hold also during mitosis (see Text S1 and figure S5). Even though this is a rough approximation of the overall mechanical behavior of cells, there are many details that must be managed to make it work, and they are all described in Text S1.

Here the Delaunay triangulation that we use as the scaffolding for the diffusion problem turns out to be useful once again: the same cell-cell links also define the set of neighbors of each cell, and therefore the global problem of computing the pairwise interactions between cells can be reduced to a single loop over all cells and the small limited number of their immediate neighbors, so that this operation has an O(N) computational complexity only – and it does not grow when we include the cost of the Delaunay triangulation [35] – instead of the O(N<sup>2</sup>) complexity of generic pairwise interactions.

#### Results

The first and most obvious result is the outstanding match of the growth curves of simulated spheroids with those of real spheroids: figure 2 shows a few stages of the growth of a simulated spheroid (a real spheroid is shown for comparison in figure 3), while figure 4 compares the growth curve of a single simulated spheroid with the growth curves of real spheroids grown in vitro. Here we see that the growth curves are very much alike, and we found that simulation runs with different parameters – in the biophysically meaningful ranges – produce very similar growth curves, in spite of structural internal changes: the growth curves are thus rather robust with respect to parameter changes.

Several experiments [37–42] have yielded many accurate measurements of oxygen and glucose concentrations and other quantities vs. spheroid radius; these values are part of the output of our simulation program as well (see figure 5 and figure 6), and a comparison with the experimental data is shown in table 1. On the whole the agreement of simulation data of single spheroids with the experimental values is quite good, and we wish to stress that this is not the result of a fit a posteriori, but rather of the a priori choice of model features and parameters. These results qualify as true predictions of the numerical model.

The necrotic core of spheroids is another important feature that is well reproduced in the simulations, and it is clearly visible in central slices of the simulated spheroid in figure 2. The simulations also provide detailed, quantitative snapshots of the necrotic core dynamics; the left column of figure 7 shows the percentage of dead cells vs. distance from the centroid of a simulated spheroid at different times. In these snapshots we can clearly observe the formation of the sharp step that marks the edge of the necrotic core.

These results indicate that the simulation program is reliable and robust and reproduces - both quantitatively and qualitatively - known experimental results. However, it yields much more than just successful comparisons: figure 8 shows two views of the spheroid microenvironment that at present would be unobtainable by other means at this level of resolution. The left panel of figure 8 is a plot of the flow of glucose in the extracellular spaces of a mature spheroid, superposed on a density plot of extracellular glucose concentration, and it shows - rather unexpectedly - that there is an outward flow of extracellular glucose from the central necrotic region. In the external, viable rim the flow is inward bound, and there is a spherical shell where the flow is stationary. The right panel of figure 8 shows the corresponding plot of cell velocities in the same central slice, and we see that the velocity vectors point outward in the viable rim, while there are wellformed vortices in the central region, and the region in-between displays distinctive chaotic motions: these three regions closely match the three regions in the left panel. The right column in figure 7 shows radial velocity vs. distance from the centroid of the simulated spheroid, and sheds some more light on the nature of this structure: as more and more cells die and the necrotic core forms, the dead cells shrink and the core contracts. The contraction of the necrotic core expels the residual glucose in the extracellular spaces and produces the observed outward flow. We found that this behavior is strongly dependent on the particular value of the diffusion coefficient and on the metabolic activity of cells. In some simulations - where we used a lower value for the effective diffusion coefficient of oxygen - we observed a similar structure with oxygen as well. We remark that in the case



Space coordinates in µm

Figure 2. Snapshots of one simulated spheroid taken at different times. As the spheroid grows, a necrotic core develops in its central region, just as it happens in real spheroids. The size of the necrotic core and of the viable cell rim match real measurements. doi:10.1371/journal.pone.0013942.g002



100 µm

Figure 3. Photograph of a spheroid grown in vitro from HeLa cells in agar. The spheroid is colored with trypan blue to mark dead cells, where the necrotic core is clearly visible. The agar contains the spheroid and helps in obtaining a better spherical shape with HeLa cells, but also stifles spheroid growth because it reduces the effective diffusion coefficients in the nourishing medium, so that it cannot be directly compared to the simulated spheroid in the second column of figure 2 (which has the same size), while it is similar to the larger spheroid in third column.

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of lactate we found no such structure, and we obtained a pH value – derived from the distribution of lactate inside the spheroid – that is very close to experimental measurements: this indicates that the discretized reaction-diffusion scheme used in the simulation program performs correctly, and that the observed flows are not algorithmic artifacts.

#### Discussion

Although the program described in this paper is based on a model of individual cells that includes only the basic cell functions,



Figure 4. Growth curve of a simulated tumor spheroid (solid line). The run parameters used in this case are listed in Text S1. The symbols denote data points taken in different in vitro experiments: squares = FSA cells (methylcholantrene-transformed mouse fibroblasts) [45]; diamonds = MCF7 cells (human breast carcinoma) [19]; circles = 9L cells (rat glioblastoma) [27].

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Model of Tumor Spheroids



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Figure 5. Concentrations in the simulated spheroid. The color coded figures on the left show the partial pressure of oxygen, the concentrations of glucose and lactate in the extracellular spaces, and the pH of the extracellular environment (high values = red, low values = blue). The corresponding plots in the right column show the average values of these quantities vs. the distance from the centroid of the tumor spheroid. The small oscillations in the plots close to the spheroid surface are due to fluctuations in the averaging procedure, because the spheroid is slightly nonspherical.

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the simulation results compare very well with experimental measurements, and give strong hints on the sources of individual spheroid variability. Moreover, the images obtained in single runs reveal unexpected and interesting correlations and an elaborate structure of the tumor microenvironment that could never be observed before. This unexpected, complex microstructure - the formation of different regions, and the flows that characterize them, along with the complex velocity field - can be discerned in the flows of the other substances, though not all of them, according to their effective diffusion coefficient and their metabolism: the figures of these flows are shown at full-size as supporting information. Thus if we suppose that, in a more complete description, there are N substances that characterize the spheroid microenvironment, and assume that the spherical shell that divides the two main regions lies in the same position for all of these substances and that their effective diffusion coefficients are uncorrelated, then 2<sup>N</sup> different spheroid structures are determined by diffusion alone. The variation of some critical parameter (e.g., a slight change in the metabolic activity due to local fluctuations in the number of dead cells, and thus a change in the effective diffusion coefficients) can potentially act as a switch and determine widely different fates for similar spheroids. This variability cannot be discerned from growth experiments: the simulations that we have performed to date indicate that the growth curve alone is not enough to distinguish between such different states, because it does not change much even when important substances, like oxygen, diffuse in markedly different ways. These different states represent different biochemical configurations of tumor microenvironment, that might exert distinct selective pressures on cells during tumor evolution.

The spheroid microstructure that is well evidenced in figure 8, and in figures S8, S9, S10, S11, S12, S13, S14, S15, S16, S17, S18, S19, S20, S21 and in Movies S1, S2, S3, shows highly correlated fluctuations that produce, e.g., islets of proliferating cells



Figure 6. Plots of the normalized average intracellular concentration of lactate (green), glucose (blue), and ATP (red). These plots have been obtained in the same simulation and at the same time step as the plots of figure 5, and each concentration is normalized to its peak value. These plots indicate that cell death in the central region is due both to the accumulation of metabolites (lactate) and to metabolic stress (starvation). doi:10.1371/journal.pone.0013942.0006

in the sea of dead cells of the core, and cell and mass flows that follow preferential channels. There is a sort of spheroid-specific self-organization of the internal structure due to these correlated fluctuations. Similar cell flows have been observed in the lab and a recent review has stressed the great significance of such findings [43]: the simulations suggest that the whole topic of cell flows and extracellular diffusion should be investigated further. On the basis of the simulation results, we also conjecture that the flow of therapeutic drugs may be diverted as well, and let some viable, proliferating tumor cells escape treatment. This means that the simulation program could eventually become an important tool to design novel treatment schedules, and possibly validate the effects of anti-tumor drugs.

Certainly the model is far from complete, and we plan to add soon several new features, like a basic model of intracellular acidity, now accounted for by a simple phenomenological parameterization, and the effects of pH and salt concentration on diffusion. However, already in its present form, we believe that this numerical model is a true testbed of biological complexity and a real virtual laboratory, and also a source of important biomedical clues.

#### Methods

The simulation program is written in ANSI C++: this programming language was a natural choice from the very start for distinct reasons:

- N C++ is an object-oriented language, and in a simulation such as this, it is very natural to define objects that have a clear-cut biological meaning;
- N at present, C++ programming is supported by a vast array of scientific libraries, and this helps reducing program development time;
- the availability of the flexible and powerful C++ library CGAL [33] that handles the computational geometry structures utilized by the program (convex hulls, Delaunay triangulations and alpha shapes);
- N the availability of powerful development tools and highly optimized compilers.

The structure of the program reflects the organization explained in the paper: a layout is shown in figure 9. The functional blocks work as follows:

#### Initialization

At start, the environmental concentrations are set at their standard levels (see Text S1), and internal variables of all cells are set at approximate standard values (see Text S1 for the cells' variables and the physical values that are hard-coded in the program). During initialization, cells are allowed to grow and proliferate freely in an environment that is held fixed. The number of cells is also kept constant, and when a mitosis occurs one of the daughter cells is discarded. In this initial phase cells can have large oscillations of their metabolic parameters, and can occasionally step in parameter regions that would normally spell death: this does not occur here. Initialization lasts until the oscillations of

#### Table 1. Comparisons with experimental parameters.

Parameter	Simulation	Experiments	<sup>7</sup> Cell type	Ref.
<sup>1</sup> Glucose uptake (kg s <sup>{ 1</sup> m <sup>{ 3</sup> )	1:44:10 { 3	5:4 { 12:6 10 { 3	Rat-T1, MR1	[37]
<sup>1</sup> Lactate release (kg s <sup>{ 1</sup> m <sup>{ 3</sup> )	1:35:10 { 3	5:4 { 9:10 { 3	Rat-T1, MR1	[37]
<sup>2</sup> pO <sub>2</sub> (mmHg)	7	0–20	Rat-T1	[37]
		0–40	MR1	[37]
		20–60	EMT6/Ro	[38]
<sup>3</sup> рН	6.7	6.6	C6, H35	[39]
		6.96-6.99	U118-MG, HTh7	[40]
⁴DpH	0.77	0.41	U118-MG	[40]
		0.49+0.08	HTh7	[40]
<sup>5</sup> Viable cell rim thickness (mm)	155	200	EMT6/Ro	[38]
		142+16	HTh7	[40]
		310+28	U118-MG	[40]
		198+27	Col12	[41]
		225+26	HT29	[41]
<sup>6</sup> Hypoxic rim thickness (mm)	98	44+52	Col12	[41]
		44+52	HT29	[41]
Cell cycle distribution (%)	G0/G1 = 57.3	G0/G1 = 58 + 4	BMG-1	[42]
	S = 21.6	S = 19+1		
	G2/M = 21.1	G2/M = 23+1		

Metabolic and histologic parameters in spheroids of approximately 500 mm diameter: comparison between a single, large simulation, carried out with the parameters listed in Text S1, and experimental data.

Notes:

<sup>1</sup>Rate of glucose uptake or lactate release per viable spheroid volume (see [37]).

 $^{2}$ Central pO<sub>2</sub> tension (experiments) or estimated in the centroid (simulations).

 $^3$ pH has been determined in the central region of the spheroids. This corresponds to a sphere of radius &100mm about the centroid of the spheroid.

<sup>4</sup>Difference between environmental pH and pH 200 mm below the spheroid surface.

<sup>5</sup>In our simulations the viable cell rim thickness corresponds to the distance between the spheroid surface and the inner shell where only 5% of the cells are still alive. Experimental values have been determined by histology.

<sup>6</sup>These values corresponds to the radius of the necrotic core.

<sup>7</sup>Cell types are as follows: Rat-T1 = T24Ha-ras-transfected Rat1 cells (Rat1 = spontaneously immortalized rat embryo fibroblasts); MR1 = myc/ T24Ha-ras-cotransfected rat embryo fibroblasts; EMT6/Ro = mouse mammary tumor cells; C6 = rat glioma cells; H35 = rat hepatoma cells; U118-MG = human glioblastoma cells; HT7 = human tyroid carcinoma cells; Col12 = moderately differentiated human colon adenocarcinoma ; HT29 = poorly differentiated human colon adenocarcinoma; BMG-1 = human glioma cells.

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metabolic parameters die out. We have determined the duration of the initialization phase observing the desynchronization of a population of initially synchronized cells: when oscillations of the relative fractions of cells in each cell-cycle phase become undetectable we estimate that cells have reached a stable state. It turns out that a simulated time of  $3^{\circ}10^{6}$ s (i.e. about 35 days of simulated time) is sufficient for initialization of cell with a period of about 20 hours. Usually the starting number of cells is quite small (normally just one cell to seed the growth of a single spheroid), and initialization executes in very short real time (a few seconds).

#### Metabolism, diffusion, transport, and growth

This part of the program solves the combined differential system of equations that describe internal cell metabolism and diffusion in the extracellular spaces (described in detail in Text S1), using the implicit Euler method. This leads to a system of nonlinear equations, that are solved in turn with a variant of the Newton-Raphson method. The functional scheme of this important part of the program is shown in figure 10. We wish to stress that although the number of variables can be quite large (more than  $10^7$  loop variables), convergence is reasonably fast, because the initial concentration values are invariably very close to the final ones.

#### Cell motion

Cell motion is also regulated by differential equations and the solution uses a strategy based on a semi-implicit method (described in detail in Text S1). Volume growth is regulated by the part that handles metabolism and diffusion, therefore it is loosely coupled to cell motion. However we have implemented an updating mechanism that effectively decouples the two parts of the program: this means that the program can use multithreading with shared memory and exploit the features of multicore processors.

#### Cellular events

This part of the program handles discrete events, like cell-cycle transitions, mitosis and cell death. In case of mitosis it also initializes the daughter cells – using the metabolic variables of the mother cell – and allocates memory for the new cells.

#### Geometry and topology of cell cluster

Geometrical and topological informations are updated here, using calls to CGAL methods [33] that compute the convex hull of the cluster of cells, the Delaunay triangulation of cell centers, and the alpha shape of the cluster – with an alpha parameter [33] equal to  $(2r_0)^2$  where  $r_0$  is the average cell radius. This part of the



Figure 7. Fraction of dead cells (left column) and average radial velocity (right column) at different times. As the spheroid grows, the necrotic core becomes increasingly well defined, and as dead cells shrink, the radial velocity changes sign and a marked inward motion characterizes the central region. doi:10.1371/journal.pone.0013942.g007



Figure 8. Two views of the microstructure of a simulated spheroid, with about 500mm diameter and 296264 cells (183893 live cells+112371 dead cells). (Left panel): flow of extracellular glucose along a central section of the tumor spheroid (yellow arrows) superposed on the plot of glucose concentration. The length of the arrows is proportional to the glucose flow intensity projected on the plane of the section. At this stage, the necrotic core is contracting as dead cells gradually shrink, and this leads to a slow outward flow of the glucose stored in the extracellular spaces in this central region. We observe that such a behavior depends on the effective diffusion coefficient of glucose, and it disappears completely when the diffusion coefficient is high enough. This also suggests that the flow of glucose and other substances, like therapeutic drugs, is strongly dependent on the biochemistry and structure of extracellular spaces, and even small changes can lead to markedly different internal spheroid morphologies. (Right panel): individual cell velocities in the simulated spheroid. This is the same central section as in the left panel, and the velocity vectors are projected on the plane of the section. The length of each vector is proportional to the projected speed. The velocities in the viable rim show a coherent outward motion, while the velocities in the necrotic core show a rather orderly inward motion, with some vortices due to local residual cell proliferation. The region in-between is somewhat chaotic and the global structure of this plot mirrors that of the glucose flow shown on the left. The supporting information includes higher-quality versions of these figures and those of other flows. doi:10.1371/journal.pone.0013942.q008



Figure 9. Functional blocks of the simulation program. Program initialization is followed by a loop that performs biochemical and biomechanical calculations. This is followed by a check of the status of individual cells – this is where we decide whether a cell advances in the cell cycle, undergoes mitosis, or dies. Next the program computes the geometry and the topology of the cell cluster, and finally it outputs intermediate statistics and results. The loop continues until a user-defined stop condition is met. Some parts of the program can proceed in parallel (like metabolism and cell motion), and we can use multithreaded code. doi:10.1371/journal.pone.0013942.g009



Figure 10. Functional blocks of the C++ method that computes metabolic and extracellular variables. This part performs a loop that computes the solution of the nonlinear equations found in the implicit Euler integration step [29] (see also Text S1). Although the number of variables can be quite large (more than 10<sup>7</sup> variables), convergence is fast, because the initial concentration values are invariably very close to the final ones. doi:10.1371/journal.pone.0013942.g010 program uses this basic information to set all relevant geometrical and topological variables in the program.

#### Summary statistics and dump on file

The last step in the loop computes several statistics and outputs them on summary files. It also writes periodically the whole configuration of cells on file for further processing.

#### Program termination

The program repeats the loop until one of the stop conditions is met: either all cells are dead, or the program executed the required number of steps. Text S1 contains additional considerations on algorithmic complexity and on measured performance (see also figure S6 and figure S7).

Additional processing to extract useful informations from the simulation data is performed with several standard tools, like Mathematica [44].

#### Supporting Information

Text S1 Including tables and additional references. Found at: doi:10.1371/journal.pone.0013942.s001 (0.89 MB DOC)

Figure S1 Sketch of the cell phases included in the simulation program. The arrow lengths suggest the relative duration of each phase. Phase G1 is divided in two parts: an initial subphase where cells are sensitive to variations in the environmental nutrient concentration (G1m) and a later subphase where cells are insensitive to deprivation of nutrients (G1p); in between these subphases there is an energetic checkpoint [1–3].

Found at: doi:10.1371/journal.pone.0013942.s002 (0.19 MB TIF)

Figure S2 Sketch of the metabolic network. Variables within circles represent molecular species and are expressed in units of concentration or mass. Non-obvious symbols are as follows (the suffixes ext and int denote, respectively, extracellular and intracellular variables): G = glucose, G6P = glucose-6-phosphate, STORE = glucose stored in the form of glycogen, AL = lactic acid, A = glutamine, ATPp = pool of ATP molecules, DNA = nuclearmass of DNA (normalized to 1 for the whole genome), mtDNA = mitochondrial DNA. Rates are represented by squares. The red circuit represents the oxygen sensor, whereas the green circuit represents the ATP sensor [1,2]. "Cell cycle checkpoints" denotes the molecular circuit of cell cycle control that has been modeled on the basis of previous studies on the dynamics of the allosteric effect [19,20]. The biological foundations of this simplified metabolic network have been given in references 1 and 2. Recent improvements with respect to our previous model include: internalization rates of glucose, glutamine and lactate are sensitive to extracellular pH and this dependence is now described by smoothed functions (see the text for details); synthesis of cellular proteins, nuclear DNA and mitochondrial DNA are now described by double-substrate Michaelis-Menten chemical reactions to take into account the dependence of protein and nucleic acid biosynthesis on glutamine (which stands phenomenologically for the wider class of aminoacids) and ATP availability.

Found at: doi:10.1371/journal.pone.0013942.s003 (0.77 MB TIF)

Figure S3 The geometry and topology of diffusion. a) Most substances are carried into and out of the cell by facilitated diffusion and there is an active mass exchange between cell and extracellular space. Each cell in the simulation program has its own extracellular space. b) Extracellular spaces are interconnected and there is a diffusion flow through the network of connections. c) The network of interconnected spaces is defined by a Delaunay

triangulation. In this 2D representation, for any red dot we can define a Voronoy cell, i.e., the set of points in the plane that are closer to the given dot than to any other dot in the set. The dual structure is the Delaunay triangulation (Voronoy cells are black and Delaunay links are green). There is a Delaunay link between any two dots only if the respective Voronoy cells touch each other, therefore we can use the Delaunay triangulation to define proximity. This enables us to set up a discretized version of diffusion between extracellular spaces, like in b). In addition to the topology of contacts between cells we also keep into account geometry:  $g_{ab}$  in part b) is a geometric factor that modulates diffusion. d) The actual simulation is in 3D: here the Delaunay triangulation of a small cluster of cell centers shows up in transparency.

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Figure S4 Pictorial representation of the interaction force between two cells. The solid curve shows qualitatively the behavior of the interaction force, while the insects depict the corresponding situations (A. cells are compressed against each other, force is repulsive; B. cells are in contact, the total force vanishes; C. cell centers are slightly apart, force is attractive because of adhesive molecules on the cells' membranes; D. cells are no longer in contact, the total force vanishes again). The inset on the right corner shows the definitions of the basic geometric variables.

Found at: doi:10.1371/journal.pone.0013942.s005 (0.36 MB TIF)

Figure S5 The geometry of mitosis.  $R_0$  is the radius of the initial cells, while  $R_1$  and  $R_2$  are the radii of the daughter cells: because of random asymmetries during mitosis, the daughter cells usually have different sizes. The program places the two daughter cells inside the region initially occupied by the mother, and the axis connecting the centers points in a random direction. This forces the two cells to push one against the other, and as cells separate the axis rotates, so that the new configuration fits the positions of neighboring cells. The cells' centers are separated by a distance which is roughly 0.4  $R_0$ .

Found at: doi:10.1371/journal.pone.0013942.s006 (0.16 MB TIF)

Figure S6 CPU time needed to simulate 1 hour, vs. the number of cells in the spheroid. In this run, the precision of the solution of the global diffusion transport and metabolism problem is fixed at 1% and the timestep is 50 s (so that  $t_{CPU}$  is actually the CPU time needed to simulate 72 elementary timesteps). The solid curve shows the fit (S.69) in Text S1.

Found at: doi:10.1371/journal.pone.0013942.s007 (0.15 MB TIF)

Figure S7 Total CPU time vs. N. This figure shows the total CPU time vs. N in the same run as figure S6. The solid curve is a simple fit with a quadratic polynomial function.

Found at: doi:10.1371/journal.pone.0013942.s008 (0.14 MB TIF)

Figure S8 Oxygen concentration and flow at simulated time = 14 days. Cells are coloured to show the oxygen concentration (color mapping, blue = low concentration, red = high concentration), while the yellow arrows show the oxygen flow (arrow length proportional to flow intensity).

Found at: doi:10.1371/journal.pone.0013942.s009 (2.06 MB TIF)

Figure S9 Oxygen concentration and flow at simulated time = 16 days. Cells are coloured to show the oxygen concentration (color mapping, blue = low concentration, red = high concentration), while the yellow arrows show the oxygen flow (arrow length proportional to flow intensity).

Found at: doi:10.1371/journal.pone.0013942.s010 (1.14 MB TIF)

Figure S10 Oxygen concentration and flow at simulated time = 17 days. Cells are coloured to show the oxygen concentra-

tion (color mapping, blue = low concentration, red = high concentration), while the yellow arrows show the oxygen flow (arrow length proportional to flow intensity).

Found at: doi:10.1371/journal.pone.0013942.s011 (4.88 MB TIF)

Figure S11 Extracellular glucose concentration and flow at simulated time = 14 days. Cells are coloured to show the concentration of extracellular glucose (color mapping, blue = low concentration, red = high concentration), while the yellow arrows show the flow of extracellular glucose (arrow length proportional to flow intensity).

Found at: doi:10.1371/journal.pone.0013942.s012 (2.05 MB TIF)

Figure S12 Extracellular glucose concentration and flow at simulated time = 15 days. Cells are coloured to show the concentration of extracellular glucose (color mapping, blue = low concentration, red = high concentration), while the yellow arrows show the flow of extracellular glucose (arrow length proportional to flow intensity).

Found at: doi:10.1371/journal.pone.0013942.s013 (2.88 MB TIF)

Figure S13 Extracellular glucose concentration and flow at simulated time = 16 days. Cells are coloured to show the concentration of extracellular glucose (color mapping, blue = low concentration, red = high concentration), while the yellow arrows show the flow of extracellular glucose (arrow length proportional to flow intensity).

Found at: doi:10.1371/journal.pone.0013942.s014 (1.15 MB TIF)

Figure S14 Extracellular glucose concentration and flow at simulated time = 17 days. Cells are coloured to show the concentration of extracellular glucose (color mapping, blue = low concentration, red = high concentration), while the yellow arrows show the flow of extracellular glucose (arrow length proportional to flow intensity).

Found at: doi:10.1371/journal.pone.0013942.s015 (4.97 MB TIF)

Figure S15 Extracellular glutamine concentration and flow at simulated time = 14 days. Cells are coloured to show the concentration of extracellular glutamine (color mapping, blue = low concentration, red = high concentration), while the yellow arrows show the flow of extracellular glutamine (arrow length proportional to flow intensity).

Found at: doi:10.1371/journal.pone.0013942.s016 (2.06 MB TIF)

Figure S16 Extracellular glutamine concentration and flow at simulated time = 17 days. Cells are coloured to show the concentration of extracellular glutamine (color mapping, blue = low concentration, red = high concentration), while the yellow arrows show the flow of extracellular glutamine (arrow length proportional to flow intensity).

Found at: doi:10.1371/journal.pone.0013942.s017 (4.94 MB TIF)

Figure S17 Extracellular glutamine concentration and flow at simulated time = 18 days. Cells are coloured to show the concentration of extracellular glutamine (color mapping, blue = low concentration, red = high concentration), while the yellow

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arrows show the flow of extracellular glutamine (arrow length proportional to flow intensity).

Found at: doi:10.1371/journal.pone.0013942.s018 (6.11 MB TIF)

Figure S18 Extracellular lactate concentration and flow at simulated time = 18 days. Cells are coloured to show the concentration of extracellular lactate (color mapping, blue = low concentration, red = high concentration), while the yellow arrows show the flow of extracellular lactate (arrow length proportional to flow intensity). Lactate always flows outward in the simulation: this is a single snapshot taken after both glucose and glutamine have developed their split flow regime.

Found at: doi:10.1371/journal.pone.0013942.s019 (6.07 MB TIF)

Figure S19 Velocity vectors projected on the plane of the slice at simulated time = 14 days. The vectors show the cells' motions in the plane of the slice (yellow arrows, arrow length proportional to flow intensity). Cells in the core perform complex looping motions, while cells in the viable rim always push outward.

Found at: doi:10.1371/journal.pone.0013942.s020 (1.62 MB TIF)

Figure S20 Velocity vectors projected on the plane of the slice at simulated time = 16 days. The vectors show the cells' motions in the plane of the slice (yellow arrows, arrow length proportional to flow intensity). Cells in the core perform complex looping motions, while cells in the viable rim always push outward.

Found at: doi:10.1371/journal.pone.0013942.s021 (3.08 MB TIF)

Figure S21 Velocity vectors projected on the plane of the slice at simulated time = 18 days. The vectors show the cells' motions in the plane of the slice (yellow arrows, arrow length proportional to flow intensity). Cells in the core perform complex looping motions, while cells in the viable rim always push outward.

Found at: doi:10.1371/journal.pone.0013942.s022 (5.00 MB TIF)

Movie S1 Movie of a central slice of a simulated tumor spheroid showing the development of the necrotic core (red = live cells, black = dead cells).

Found at: doi:10.1371/journal.pone.0013942.s023 (5.99 MB MOV)

Movie S2 Movie of a central slice of a simulated tumor spheroid showing the flow of extracellular glucose (same coding as figures S11, S12, S13, S14).

Found at: doi:10.1371/journal.pone.0013942.s024 (21.18 MB MOV)

Movie S3 Movie of a central slice of a simulated tumor spheroid showing the map of projected cell velocities (same coding as figures S19, S20, S21).

Found at: doi:10.1371/journal.pone.0013942.s025 (28.40 MB MOV)

#### Author Contributions

Conceived and designed the experiments: EM RC. Performed the experiments: EM RC. Analyzed the data: EM RC. Contributed reagents/materials/analysis tools: EM RC. Wrote the paper: EM RC.

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## Supporting text for

## Emergent Properties of Tumor Microenvironment in a Real-life Model of Multicell Tumor Spheroids

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## Introduction

This supplementary text contains a systematic overview of the mathematical and computational details of the model introduced in the main text, and is divided in six parts. We start with a review of the single-cell model which summarizes the internal biochemical machinery included in the simulated cells: although this has already been described before [1,2], here we include some recent improvements and in the second part we summarize the setup of the numerical method for diffusion and transport processes. This is followed by a review of the dynamical interactions whereby cells interact mechanically with one another. Finally we list all the parameters used in the simulation program and include broad considerations on the computational complexity of the model. The last section contains a selection of images obtained in the simulations.

## A. Brief review of the single-cell model

The cells in the simulation program behave as automata that obey mostly deterministic rules, although they also include some circumscribed but important randomness. Each cell has many individual status variables that are listed in table S1, and steps through well-defined phases before dividing. After cell division the old cell is erased from the program's cell list, the daughter cells are added to the list and each of them starts a new independent cell cycle. The cell cycle is sketched in figure S1.

We approximate cells as soft spheres of radius r, surrounded by a thin layer that represents the extracellular space. The extracellular space is essential for a proper description of diffusion of molecules like glucose that are transported across the cell membranes by facilitated diffusion, and the variables that characterize the extracellular space around each cell are also included in the list of status variables.

## The biochemistry of cells

The biochemical model of an isolated tumor cell developed in references [1] and [2] is outlined in figure S2 and comprises the following parts:

- 1. Glucose uptake and retention by cells
- 2. Glucose-6-phosphate utilization by cells
- 3. Glycogen storage
- 4. Glutamine (aminoacid) uptake
- 5. Oxygen uptake
- 6. ATP production and consumption
- 7. Lactate production
- 8. DNA synthesis
- 9. Protein synthesis
- 10. Mitochondrial proliferation
- 11. Dynamics of cell volume
- 12. Checkpoints and phase changes
- 13. Corrections due to internal and environmental pH
- 14. Other corrections

Now we describe in detail each item in the list.

Glucose uptake and retention by cells

The equations that regulate the metabolism and transport of glucose are

$$\frac{dm_{G,C}}{dt} = M_G \Big[ m_{G,C}(t) \Big] + T_G \Big[ m_{G,C}(t), m_{G,C}(t) \Big]$$
(5.1a)

$$\frac{dm_{G,c}}{dt} = D_G \sum_{\langle b \rangle} \left( \frac{m_{G,b}(t)}{V_b(t)} - \frac{m_{G,c}(t)}{V_c(t)} \right) g_{bc} - T_G \left[ m_{G,c}(t), m_{G,C}(t) \right]$$
(S.1b)

where  $m_{G,C}$  and  $m_{G,c}$  denote respectively the mass of glucose inside the cell (C) and inside the surrounding extracellular volume (c).

The metabolic ( $M_G$ ) and the transport ( $T_G$ ) functions are respectively

$$M_{G}(m_{G,C}) = -\frac{v_{\max,2}m_{G,C}^{2}}{\left(V_{C}K_{2} + m_{G,C}\right)\left(V_{C}K_{a} + m_{G,C}\right)} - \frac{v_{\max,22}m_{G,C}^{2}}{\left(V_{C}K_{22} + m_{G,C}\right)\left(V_{C}K_{a} + m_{G,C}\right)}$$
(5.2*a*)

$$T_{G}(m_{G,C}, m_{G,c}) = a2c \cdot \frac{v_{\max,1}m_{G,c}}{V_{c}K_{1} + m_{G,c}} - c2a \frac{v_{\max,1}m_{G,C}}{V_{c}K_{1} + m_{G,C}}$$
(S.2b)

where

 $v_{\text{max},1} = VMAX1 \cdot h \cdot S$ 

*h, a2c, c2a* are correction factors that shall be specified later on, *S* is the cell surface,  $V_c$  is the cell volume,  $V_c$  is the extracellular volume, and finally the values  $v_{max,2} = VMAX2$ ;  $v_{max,22} = VMAX22$ ;  $K_1$ ,  $K_2$ ,  $K_{22}$ , are cellular parameters (see below).

The metabolic function ( $M_G$ ) in equation (S.1a) corresponds to hexokinase and glucokinase activity<sup>1</sup>.

The first term on the rhs of equation (S.1b) describes diffusion on the disordered lattice of extracellular volumes; here  $D_G$  is the glucose diffusion coefficient either in the extracellular space or in the enviroment (water) (see table S5), the  $g_{bc}$ 's are numerical coefficients that belong to the discretization of the diffusion problem on the cell network (see section B), and the triangular bracket  $\langle b \rangle$ denotes the set of the adjacent extracellular spaces. This diffusion part is necessary for the consistency of the whole procedure: diffusion in the cell cluster and transport into and out of the cell are inescapably related (see section B and ref. 4).

## Glucose-6-phosphate (G6P) utilization by cells

The equation that regulates the G6P content of the cell is

$$\frac{dm_{G6P,C}}{dt} = -M_G(m_{G,C}) - (g_1 + g_2 + g_3)$$
(5.3)

where  $m_{GGP,C}$  is the mass of G6P inside cell C, and the rates  $g_i$  denote different paths of G6P utilization (see also figure S2 and reference [1])

$$g_1 = \left(t_{pH} \cdot coeffg1\right) m_{(G6P),C}$$
(S.4a)

$$g_2 = (coeffg2:SensO2)m_{(G6P),C}$$
(S.4b)

$$g_3 = (coeffg3)m_{(G6P),C}$$
(5.4c)

The underlying utilization model is straighforward, all the rates  $g_i$  are simply proportional to G6P content of the cell, as this fits quite well the existing experimental data [1].

Moreover  $t_{pH}$  and SensO2 are correction terms (see below) and coeffg1, coeffg2, coeffg3 are cellular parameters (see below).

The metabolic function ( $M_G$ ) has been defined above for glucose, and corresponds to glucose transformation into G6P.

## Glycogen storage

Glucose is partly stored as glycogen. Its production/consumption, follows the equation

$$\frac{d}{dt}Store = (g_3 + p_{11}) - (r_1 + r_2 + r_3)$$
(5.5)

where Store is the glycogen mass inside the cell, the rates  $r_i$  correspond to different paths of Store utilization, the rate  $p_{11}$  corresponds to glycogen

production via the glutamine pathway (defined below), and finally the rate  $g_3$  is associated to the glycogen production path from G6P conversion

$$r_{1} = \left(t_{pH} \cdot coeffr1\right) \frac{Store}{V_{C} \cdot Kmc + Store}$$
(S.6a)

$$r_2 = g_3 \cdot SensO2 \frac{Store}{V_C \cdot Kmc + Store}$$
(5.6b)

$$r_{3} = SensO2 \cdot SensATP \cdot \frac{1}{30} \cdot \frac{PM_{G}}{PM_{ATP}} \cdot (ATP_{St} - ATPOx) \frac{Store}{V_{C} \cdot Kmc + Store}$$
(S.6c)

$$p_{11} = t_{p11} \left( t_{pH} \cdot coeffr1 \right) \cdot \frac{m_{A,C}}{V_C \cdot Kmd + m_{A,C}} \cdot \left( 1 - SensATP \right)$$
(S.6d)

Notice that these rates are modeled with simple Michaelis-Menten functions, where the substrate is the *Store* itself (for consumption) or glutamine (for production).

In addition to the symbols already defined in the previous sections, there are also other correction factors,  $t_{p11}$ , SensATP, and SensO2 (defined later). There is also a couple of additional parameters, Kmd and Kmc, and two molecular weights,  $PM_G$ and  $PM_{ATP}$  (all quantities identified by symbols like  $PM_X$  are molecular weights): they are all listed in table S2. There are also two quantities related to ATP,  $ATP_{St}$ and ATPOx, which are defined in the section on ATP.

## Glutamine (aminoacid) uptake

The equations that regulate the metabolism and transport of glutamine are

$$\frac{dm_{A,C}}{dt} = M_A \Big[ m_{A,C}(t) \Big] + T_A \Big[ m_{A,c}(t), m_{A,C}(t) \Big]$$
(5.7a)

$$\frac{dm_{A,c}}{dt} = D_A \sum_{\langle b \rangle} \left( \frac{m_{A,b}(t)}{V_b(t)} - \frac{m_{A,c}(t)}{V_c(t)} \right) g_{bc} - T_A \left[ m_{A,c}(t), m_{A,C}(t) \right]$$
(S.7c)

where  $m_{A,C}$  and  $m_{A,c}$  denote respectively the mass of glutamine inside the cell (C) and inside the surrounding extracellular volume (c). As for glucose, equation (S.7c) includes a diffusion term and  $D_A$  is the diffusion constant of glutamine either in the extracellular space or in the environment (water) (see table S5).

The metabolic  $(M_A)$  and the transport  $(T_A)$  functions are respectively

$$M_{A}(m_{A,C}) = -(p_{11} + p_{22})$$
  
$$-\frac{VMAXP_{A}ATPp \cdot m_{A,C}}{V_{C}^{2}Kmp + ATPp \cdot m_{A,C}} - \frac{VMAXDNA_{A}ATPp \cdot m_{A,C}}{V_{C}^{2}KmDNA + ATPp \cdot m_{A,C}} - \frac{VMAXM_{A} \cdot ATPp \cdot m_{A,C}}{V_{C}^{2}KmM + ATPp \cdot m_{A,C}}$$
  
(5.8a)

$$T_{A}(m_{A,C}, m_{A,c}) = a2cA \cdot \frac{v_{\max,A}m_{A,c}}{V_{c}KmA + m_{A,c}} - c2aA \cdot \frac{v_{\max,A}m_{A,C}}{V_{c}KmA + m_{A,C}}$$
(5.8b)

where  $v_{\text{max},A} = VMAXA \cdot S$ , and

$$p_{22} = 3 \cdot SensO2 \cdot SensATP \cdot \frac{1}{30} \cdot \frac{PM_G}{PM_{ATP}} \cdot (ATP_{St} - ATPOx) \frac{m_{A,C}}{V_C \cdot Kmd + m_{A,C}}$$
(5.9)

is the rate of glutamine consumption for ATP production (see below), and the three double-substrate Michaelis-Menten rates correspond to protein production, DNA synthesis, and glutamine consumption for mitochondrial proliferation. In these formulas there are two additional correction factors, *a2cA* and *c2aA*, defined below, and several parameters, *VMAXA*, *VMAXP*<sub>A</sub>, *VMAXDNA*<sub>A</sub>, *VMAXM*<sub>A</sub>, *Kmp*, *KmDNA*, *KmM*, all listed below.

## Oxygen uptake

The equation for oxygen metabolism and diffusion is

$$\frac{dm_{O2,C}}{dt} = D_{O2} \sum_{\langle B \rangle} \left( \frac{m_{O2,B}(t)}{V_B(t)} - \frac{m_{O2,C}(t)}{V_C(t)} \right) g_{BC} + M_{O2} \left[ m_{O2,C}(t) \right]$$
(5.10)

where  $D_{O2}$  is the oxygen diffusion coefficient either in water or in the cluster of cells (see table S5), and the metabolic function is

$$M_{O2}(m_{O2,C}) = -6 \frac{PM_{O2}}{PM_G} (g_2 + r_2 + r_3 + p_{22})$$
(5.11)

There is no extracellular quantity here, as oxygen diffuses normally over the cell and there is no specific transport process.

## ATP production and consumption

The differential equation for ATP is

$$\frac{d}{dt}ATPp = ATPtot \tag{5.12}$$

where *ATPp* is the total ATP mass in the cell, and *ATPtot* is the total variation rate. *ATPtot* is a sum of several contributions:

- *ATPOx* (ATP production rate from oxidative phosphorylation)
- *ATPNOx* (ATP production rate from anaerobic glycolysis)
- *ATP2* (ATP production rate from glycogen (Store))
- ATP3 (ATP production rate from aminoacids (glutamine) )
- ConsATP (ATP consumption rate due to gluconeogenesis)
- ConsATP<sub>1</sub> (parameterization of the ATP consumption proportional to cellular volume)
- ConsATP<sub>2</sub> (ATP consumption rate for protein synthesis)
- *ConsATP*<sub>3</sub> (ATP consumption rate for DNA synthesis)
- *ConsATP*<sub>5</sub> (ATP consumption rate for mitochondrial proliferation)

which are defined by the equations

$$ATPOx = 30 \frac{PM_{ATP}}{PM_{G}} (g_{2} + r_{2})$$
(5.13.a)

$$ATPNOx = 2\frac{PM_{ATP}}{PM_{G}}(g_{1}+r_{1})$$
(5.13.b)

$$ATP2 = 30 \frac{PM_{ATP}}{PM_G} r_3$$
(5.13.c)

$$ATP3 = 28 \frac{PM_{ATP}}{PM_{G}} p_{22}$$
(5.13.d)

$$ConsATP = 2 \frac{PM_{ATP}}{PM_G} p_{11}$$
(S. 13.e)

 $ConsATP_1 = v_{work}C1 \cdot ATPp \tag{S.13.f}$ 

$$ConsATP_{2} = \frac{VMAXP_{ATP}ATPp \cdot m_{A,C}}{V_{C}^{2}Kmp + ATPp \cdot m_{A,C}}$$
(S.13.g)

$$ConsATP_{3} = \frac{VMAXDNA_{ATP}ATPp \cdot m_{A,C}}{V_{C}^{2}KmDNA + ATPp \cdot m_{A,C}}$$
(S.13.h)

$$ConsATP_{5} = \frac{VMAXM_{ATP} \cdot ATPp \cdot m_{A,C}}{V_{C}^{2}KmM + ATPp \cdot m_{A,C}}$$
(5.13.i)

$$ATPtot = ATPOx + ATPNox + ATP2 + ATP3 -ConsATP - ConsATP_1 - ConsATP_2 - ConsATP_3 - ConsATP_5$$
(5.13.j)

In these equations there are a few additional parameters,  $v_{work}$ , C1,  $VMAXP_{ATP}$ ,  $VMAXDNA_{ATP}$ , and  $VMAXM_{ATP}$ , which are specified below.

Here we remark that some of these terms are representative of large classes of processes, in particular  $ConsATP_1$  represents all those processes that utilize ATP and are extensively distributed over the cellular volume<sup>1</sup>.

## Lactate production

The equations that regulate the production and transport of lactate are

$$\frac{dm_{AcL,C}}{dt} = M_{AcL} \left[ m_{AcL,C}(t) \right] + T_{AcL} \left[ m_{AcL,c}(t), m_{AcL,C}(t) \right]$$
(S. 14a)

$$\frac{dm_{AcL,c}}{dt} = D_{AcL} \sum_{\langle b \rangle} \left( \frac{m_{AcL,b}(t)}{V_b(t)} - \frac{m_{AcL,c}(t)}{V_c(t)} \right) g_{bc} - T_{AcL} \left[ m_{AcL,c}(t), m_{AcL,C}(t) \right]$$
(S. 14.b)

where  $D_{AcL}$  is the lactate diffusion coefficient either in the extracellular space or in the enviroment (water) (see table S5), and

$$M_{AcL}(m_{AcL,C}) = 2g_1 + 2r_1$$

$$T_{AcL}(m_{AcL,C}, m_{AcL,c}) = a2cAcL \cdot \frac{v_{\max,AcL}m_{AcL,c}}{V_c \cdot KmAcL + m_{AcL,c}} - c2aAcL \cdot \frac{v_{\max,AcL}m_{AcL,C}}{V_C \cdot KmAcL + m_{AcL,C}}$$
(S. 15a)
$$(S. 15b)$$

where  $v_{\max,AcL} = VMAXAcL$ 's, and, as before, *a2cAcL* and *c2aAcL* are correction factors, while VMAXAcL and KmAcL are cellular parameters.

## DNA synthesis

The equation for DNA synthesis is

<sup>&</sup>lt;sup>1</sup> Notice also that a term *ConsATP*<sub>4</sub> is missing: the name is reserved to a class of processes not yet included in the program.

$$\frac{dDNA}{dt} = v_{DNA}$$
(5.16)

where DNA is the fraction of a complete DNA molecule, and  $v_{DNA}$  is the rate of DNA synthesis from a double-substrate Michaelis-Menten term:

$$v_{DNA} = \frac{VMAXDNA \cdot m_{ATP} \cdot m_{A}}{V_{C}^{2} KmDNA + m_{ATP} \cdot m_{A}}$$
(5.17)

Here the glutamine mass  $(m_A)$  represents the contribution of the other aminoacids as well.

## Protein synthesis

The equation for protein production is similar to (S.16), with the difference that now it represents a whole class of substances (all the proteins):

$$\frac{dm_p}{dt} = v_p \tag{S.18}$$

where  $m_p$  is the total protein mass in the cell, and  $v_p$  is the rate of protein production from a double-substrate Michaelis-Menten term:

$$v_p = \frac{VMAXP \cdot ATPp \cdot m_A}{V_C^2 Kmp + ATPp \cdot m_A}$$
(5.19)

The protein mass determines also the amounts (masses) of some specific proteins, pRb (retinoblastoma protein) and three cyclins (D, E, X = A+B):

pRb\_fraction = 1.5e-2 times m<sub>p</sub>
cyclinD\_fraction = 10.e-3 times m<sub>p</sub>
cyclinE\_fraction = 3.e-3 times m<sub>p</sub>
cyclinX\_fraction = 14.e-3 times m<sub>p</sub>

The amount of pRb that is carried over to daughter cells in mitosis is important because this indirectly determines the duration of the G1-phase, and we model the fluctuating duration of the G1-phase assuming that pRb is bound to parts of the nucleus that are partitioned binomially at mitosis. This modeling requires one additional cell parameter, NUCLEAR\_OBJ, which is the (fictitious) number of parts to which pRb is bound.

## Mitochondrial proliferation

The number of mitochondria is described by the continuous variable *M*, which follows the differential equation

$$\frac{dM}{dt} = v_M \tag{5.20}$$

where  $v_{M}$  is the rate of mitochondrial proliferation and is associated to a doublesubstrate Michaelis-Menten term which is proportional to the production of mitochondrial DNA:

$$v_{M} = \frac{VMAXM \cdot ATPp \cdot m_{A,C}}{V_{C}^{2}KmM + ATPp \cdot m_{A,C}}$$
(S.21)

where VMAXM is yet another cellular parameter.

Mitochondria are partitioned randomly between the daughter cells at mitosis. Experimental observations indicate that this partitioning follows a binomial distribution [5,6], however, in order to better reproduce the observed desynchronization of disperse cell populations [7], mitochodria are clustered in small groups - a noteworthy fact that is experimentally observed [8] - and it is these groups that are actually shared between cells; the clustering factor *ClusteringFactor* belongs to the list of cell parameters as well.

## Dynamics of cell volume

Cellular volume is a notoriously complicated cellular variable [9,10], and in particular it is related to osmotic pressure, and thus to the concentration of many substances inside cells. Here we take the total ATP mass as representative of this vast array of substances, and we assume that the cellular volume is partly determined by the total ATP mass. In addition, there are fixed volume contributions from the cell nucleus and from the organelles. The organelles themselves are variable in number, and again, we take the number of mitochondria as representative of the whole class of cellular organelles. These considerations yield the following - extremely simplified - formula for the volume of living cells:

$$V_{C} = V_{\min}(1 + DNA) + C2 \cdot Mit + C1 \cdot ATPp$$
(5.22)

where DNA is the fraction of synthetized DNA molecule as in eq. (S.16), while C1, C2 and  $V_{min}$  are cellular parameters. Since  $V_{min}$  is the nuclear volume, eq. (S.22) takes into account the effective doubling of nuclear volume during the S-phase.

In the program we also assume that a cell can become apoptotic - for reasons explained below, in the section on cell-phase transitions - and after death the volume gradually shrinks [11] according to the equation

$$\frac{dV_c}{dt} = -DVap \cdot V_c \tag{S.23}$$

which has the exact solution

$$V_c = V_0 \exp(-DVap t) \tag{5.24}$$

where *DVap* is yet another cellular parameter, and  $V_0$  is the cell's volume at death. This shrinking continues until  $V_c = V_{min}$  and then stops: thus the cluster of cells eventually contains a sizeble volume of dead cells' residues.

## Corrections due to internal and environmental pH

The equations above include several correction factors that modify the enzyme activity according to the internal and the extracellular pH

## $pH_c = pH$ inside cell C

 $pH_c = pH$  in the extracellular space

We plan to include in the program a rather complete model of cellular acidity, but at the moment the internal pH is fixed and the external pH is computed from a simple model of the buffering capacity of the standard nutrient solutions [2]. If [AcL] is the lactate concentration, then the pH of the environment and of the extracellular spaces is

$$pH = 7.5443 - \frac{[AcL]}{BufCapEnv}$$
(S.25)

where<sup>2,12</sup> BufCapEnv = 
$$0.19953 \text{ kg/m}^3$$
.

The pH value is used to compute the following correction factors [2,13-16]:

 $t_{pH}$  = correction of metabolic rates (depends on internal pH, and thus at the moment it is fixed)

 $t_{p11}$  = fine tuning of  $p_{11}$  (depends on internal pH, and thus at the moment it is fixed)

a2c = correction to transport of glucose from extracellullar space into cell c2a = correction to transport of glucose from cell to extracellullar space a2cA = correction to transport of glutamine from extracellullar space into cell c2aA = correction to transport of glutamine from cell to extracellullar space a2cAcL = correction to transport of lactate from extracellullar space into cell c2aAcL = correction to transport of lactate from cell to extracellullar space The defining equations are:

$$t_{pH} = \frac{1}{2} \Big[ 1 + \tanh(tpH_slope \cdot pH_c - tpH_thr) \Big]$$
(S.26a)

$$t_{p11} = \frac{1}{2} \Big[ 1 + \tanh(tp11\_slope \cdot pH_c - tp11\_thr) \Big]$$
(5.26b)

$$a2c = \frac{1}{2} \left[ 1 + \tanh\left(a2c\_slope \cdot pH_c - a2c\_thr\right) \right]$$
(5.26c)

$$c2a = \frac{1}{2} \left[ 1 + \tanh\left(c2a\_slope \cdot pH_c - c2a\_thr\right) \right]$$
(S.26d)

$$a2cA = \frac{1}{2} \Big[ 1 + \tanh\left(a2cA\_slope \cdot pH_c - a2cA\_thr\right) \Big]$$
(5.26e)

$$c2aA = \frac{1}{2} \Big[ 1 + \tanh(c2aA\_slope \cdot pH_c - c2aA\_thr) \Big]$$
(S.26f)

$$a2cAcL = 2 - \tanh(a2cAcL\_slope \cdot pH_c - a2cAcL\_thr)$$
(S.26g)

$$c2aAcL = 2 - \tanh(c2aAcL\_slope \cdot pH_c - c2aAcL\_thr)$$
(S.26h)

These equations are smoothed versions of piecewise linear approximations of experimental data [2,13-16].

## Other corrections

There are a few additional correction factors related to oxygen and ATP [1,2]:

h = correction to glucose transport due to oxygen concentration
 SensO2 = correction to glucose metabolism due to oxygen concentration
 SensATP = this factor is used to switch ATP production between oxidative
 phosphorylation and anaerobic glycolysis

The equations that define these correction factors are [1,2]:

$$h = 0.5 \left[ 1.3 \left( 1 - \frac{m_{O2,C}}{V_C O2 st} \right) + 1 \right] \cdot \left\{ 1 + \tanh \left[ 100 \left( 1 - \frac{m_{O2,C}}{V_C O2 st} \right) \right] \right\}$$
  
+0.5  $\left\{ 1 - \tanh \left[ 100 \left( 1 - \frac{m_{O2,C}}{V_C O2 st} \right) \right] \right\}$  (5.27)

$$SensO2 = \frac{m_{O2,C}}{V_C K_{O2} + m_{O2,C}}$$
(5.28)

$$SensATP = \frac{1}{2} \left[ 1 - \tanh\left(100 \cdot \frac{\left(ATPOx - ATP_{St}\right)}{ATP_{St}}\right) \right] = \frac{1}{2} \left\{ 1 - \tanh\left[100 \cdot \left(\frac{ATPOx}{ATP_{St}} - 1\right)\right] \right\}$$
(5.29)

The equations for h and SensATP are smoothed versions of piecewise linear approximations, and  $ATP_{St}$  is a standard rate (this is a cellular parameter as well), that acts as a threshold that separates the two regimes of oxidative phosphorylation and anaerobic glycolysis. There are a two more parameters in these definitions, O2St, the standard oxygen concentration, and  $K_{O2}$ . ATPOx is the rate of ATP production from oxidative phosphorylation (defined above, in the section on ATP). The hyperbolic tangents that occur in the formulas do not have any special meaning, they are only practical approximations of smoothed step functions.

## Cell phase transitions

The transitions from one cell phase to the next are regulated by a complex biochemical machinery that we include only in small part, and as it interacts in deep ways with cell metabolism.

Each cell in the simulation is in a certain phase of the cell's cycle, and there are specific rules that determine the transition from one phase to the next. We define the following phases (sketched in fig. S1):

- 1. G1m phase: G1 phase prior to the commitment to enter S phase (variable duration);
- 2. G1p phase: G1 phase after the commitment to enter S phase (variable duration);
- 3. S phase (variable duration);
- 4. G2 phase (variable duration);
- 5. M phase (fixed duration);
- 6. Death (indefinite duration);

The dynamics of the G1m-G1p and of the G1p-S transitions is regulated by the pRb protein: pRb forms a complex with an enzyme E and carries 16 putative phosphorylation sites [17,18]. In this context, the enzymatic reaction catalyzed by E is a phenomenological model of the reactions controlled by the transcription factor E2F that associates with pRb in real cells.

The pRB-E complex is synthesized in the G2 phase and it is partitioned at random between the daughter cells at mitosis. At the beginning of the G1m phase the cyclin D protein (CycD) is expressed and it phosphorylates the pRB-E complex upon rapid association with specific cyclin-dependent kinases [17,18]. The phosphorylation event is assumed to occur following the reversible bimolecular interaction between cyclin D and the pRb complex: a detailed study has shown that the precise mechanism of pRb phosphorylation is irrelevant with respect to the dynamics of the system's response [17]. Upon partial phosphorylation of the pRb protein, a fraction of E molecules are released and catalyze a reaction whereby a substrate S is converted into a product P. When the concentration of S falls below the threshold Thresh\_S\_start the expression of the cyclin E protein (CycE) is activated and this marks the progression of the cell from the G1m phase to the following G1p-phase. CycE participates to pRb phosphorylation upon instantaneous association with specific cyclin-dependent kinases and leads to the hyperphosphorylation of the pRb with full E detachment. When the concentration of S falls below the threshold Thresh\_S\_stop the cell progresses to the S-phase. At the end of the G1m- and G1p-phases, the cyclins CycD and CycE are destroyed by a proteolytic degradation that is assumed to occur faster than the considered reaction kinetics. The nature of pRb phosphorylation naturally

introduces a threshold and a time delay in the reaction catalyzed by E with characteristic times comparable with the cell progression through the various cell cycle phases [19,20].

The program computes first the attachment probabability

$$p = \frac{1}{2N[A]_0} \left\{ \left( N[A]_0 + [B]_0 + \frac{k_-}{k_+} \right) - \sqrt{\left( N[A]_0 + [B]_0 + \frac{k_-}{k_+} \right)^2 - 4N[A]_0[B]_0} \right\}$$
(5.30)

and then the fraction and the number of molecules with at least *k* occupied phosphorylation sites:

$$P_{k} = P(\text{occupied sites} \ge k) = \sum_{l=k}^{N} \binom{N}{l} p^{l} (1-p)^{N-l}$$
(5.31)

$$(N_A V[A]) \cdot P_k = N_A V[A] \sum_{l=k}^{N} {N \choose l} p^l (1-p)^{N-l}$$
 (5.32)

where [A] is the concentration of the pRb protein, and [B] is the concentration of the phosphorylating complex, which - as explained above - we identify with the sum of the D and E cyclins (that we assume to be phosphorylated, in a cytoplasm with plenty of ATP). Here N and k are cell parameters (although in this case we do not expect them to change in different cell lines - see below).

Finally *pRb\_ONOFFratio* is the ratio of the detachment-attachment rates (of the mediated phosphorylation process of *pRb*):

$$\frac{k_{-}}{k_{+}} = pRb_ONOFFratio$$

We remark that all the quantities related to pRb phosphorylation process depend directly on pRb concentration, and thus on protein mass, and they do not require additional differential equations.

The concentration of the enzyme E (E2F) released by the configurational switch of the pRb protein is given by

$$[E]_{0} = [A]_{0} \sum_{l=k}^{N} {N \choose l} p^{l} (1-p)^{N-l}$$
(5.33)

and this determines the variation of the concentration of substrate S which leads to the threshold mechanism for the G1m-G1p commitment point

$$\frac{d[S]}{dt} \approx -\frac{k3MM \cdot [E]_0[S]}{KmMM + [S]}$$
(5.34)

(where *k3MM* and *kmMM* are two additional cell parameters). The cell crosses the G1m-G1p commitment point when the concentration of the substrate S reaches a low value *Thresh\_S\_start*. A still lower value *Thresh\_S\_stop* corresponds to the G1-S checkpoint transition.

DNA synthesis takes place during the S-phase only, and the program contains a variable that takes care of synthesized DNA (measured in units of complete molecule). When this counter equals 1 the S-phase ends and the cell enters phase G2.

The G2-phase ends when there is a sufficient amount of cyclin A + cyclin B (at least CycXThr). Finally the M-phase has a fixed duration (M\_T\_MEAN).

In the program we also define a special cell phase which corresponds to death. We introduce an average death rate due to metabolites in the environment:

$$death \ rate = a_R \ [AcL]_C \tag{S.35}$$

where  $[AcL]_c$  is the lactate concentration inside the cell, so that the probability that a given cell survives during a timestep  $\Delta t$  is

$$\exp\left(-a_R \left[AcL\right]_C \cdot \Delta t\right) \tag{5.36}$$

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(this is the just the conventional expression for the survival probability in toxicology<sup>21</sup>).

A cell can also die if there are not enough nutrients: we check for this condition at the start of each phase, and if the ATP content of the cell is less than a minimum mass *ATPmin*, the cell dies. The value *ATPmin* is calculated assuming that any residual cytoplasmic volume - which in our minimal model of the cell's volume depends on the total ATP content of the cell - must be at least as large as the total mitocondrial volume, then from the volume formula (S.22), we find

$$ATPmin = \frac{C2 \cdot Mit}{C1}$$
(5.37)

# B. Combined numerical solution of the diffusion and transport problem

In the previous section we have introduced several non-linear differential equations: these equations are used to simultaneously model both internal cell metabolism, and transport and diffusion. The equations have a rather conventional structure, as it often happens in biology, where the models are often concerned with enzyme action in the synthesis and destruction of molecular species, and in the regulation of transport of molecules into and out of the cell or cellular compartments.

The modeling of transport and diffusion is achieved by discretization of the diffusion problem: the cells' centers provide a natural discretization scaffolding and we use a Delaunay triangulation to define proximity relations [22-28]. Figure S3 shows schematically the cells and their extracellular spaces, together with two simple examples of triangulation, both in 2D and in 3D. The molecular species in the program can be grouped in three different classes. The first class includes substances that are neither transported outside the cell nor diffuse to neighboring cells, and obey differential equations with the generic form

$$\frac{dm_c}{dt} = M\left(m_c\right) \tag{S.38}$$

where  $m_c$  is the mass of the substance in cell *C*, and  $M(m_c)$  is a function that depends on the current state of the cell (e.g., in the case of stored glycogen, this function is implicitly defined by equation (S.5) and is equal to  $(g_3 + p_{11}) - (r_1 + r_2 + r_3)$ ). In this case, the algorithmically stable solution derived from the implicit Euler method is

$$m_C^{n+1} = m_C^n + \Delta t \cdot M\left(m_C^{n+1}\right) \tag{5.39}$$

where  $m_{C}^{n}$  is the mass of substance in the cell at the *n*-th timestep.

The second class of substances behaves like oxygen, which is metabolized by the cell, but is free to diffuse from cell to cell, and obeys equations with the generic form

$$\frac{dm_C}{dt} = M\left(m_C\right) + D\sum_{\langle B \rangle} \left(\frac{m_B}{V_B} - \frac{m_C}{V_C}\right) g_{BC}$$
(5.40)

where in addition to the metabolic term there is a diffusion term as well (and the sum runs over all adjacent cells B - here adjacency is defined by the Delaunay triangulation). Now the algorithmically stable solution derived from the implicit Euler method is equivalent to the Backward Differentiation Method method for partial differential equations [29], and writes

$$m_{C}^{n+1} = m_{C}^{n} + \Delta t \left[ M\left(m_{C}^{n+1}\right) + D \sum_{\langle B \rangle} \left( \frac{m_{B}^{n+1}}{V_{B}} - \frac{m_{C}^{n+1}}{V_{C}} \right) g_{BC} \right]$$
(5.41)

The last class of substances is like glucose, which is transported across the cell membrane by facilitated diffusion. This case requires two equations per cell

$$\frac{dm_c}{dt} = M(m_c) + T(m_c, m_c)$$
(S.42a)

$$\frac{dm_c}{dt} = -T\left(m_c, m_c\right) + D\sum_{\langle b \rangle} \left(\frac{m_b}{V_b} - \frac{m_c}{V_c}\right) g_{bc}$$
(5.42b)

where the first equation describes internal metabolism and transport across the cell membrane into (or from) extracellular space, while the second equation describes the mass balance in the extracellular space, which exchanges mass with the cell and is linked to the rest of the system by diffusion (the extracellular space corresponding to cell *C* is denoted by a lowercase *c*, and the sum num runs over the adjacent extracellular spaces *b*). Now the algorithmically stable solution derived from the implicit Euler method produces the pair of equations

$$m_{C}^{n+1} = m_{C}^{n} + \Delta t \left[ M \left( m_{C}^{n+1} \right) + T \left( m_{C}^{n+1}, m_{C}^{n+1} \right) \right]$$
(S.43a)

$$m_{c}^{n+1} = m_{c}^{n} + \Delta t \left[ -T\left(m_{c}^{n+1}, m_{C}^{n+1}\right) + D \sum_{\langle b \rangle} \left(\frac{m_{b}^{n+1}}{V_{b}} - \frac{m_{c}^{n+1}}{V_{c}}\right) g_{bc} \right]$$
(5.43b)

Finally, taking into account all substances, we are left with a large set of nonlinear equations like (S.39), (S.41), (S.43a) and (S.43b). These equations are solved iteratively with Newton-Raphson steps [30]; these steps usually converge, but may occasionally fail if the mass is very close to zero, and in this case they are substituted by steps performed with the secant method [30]. The convergence condition is that for all masses the inequality

$$\left|m^{n+1} - m^n\right| < \varepsilon \cdot \frac{m^{n+1} + m^n}{2} \tag{S.44}$$

holds, where  $\varepsilon$  is a given precision, i.e., the difference between two successive values in the iteration must be less than a given fraction of the estimated mass.

## C. Biomechanical modeling

Mechanical interactions between cells are - if possible - even more complex than the biochemical network described above, and we know that cells are viscoelastic structures which are actively shaped by the cytoskeleton [31].

The simulation work faces the following difficulties:

 existing theoretical models capture only part of the actual features of cells and are not really applicable to the complex reality of cells in clusters (e.g., reference [32] describes an interesting model of erythrocytes, which are however both simpler and different from other human cells; the model is also inapplicable here, because it would lead to an exceedingly high computational complexity);
- many measurements are inapplicable to the simulation of cells (pipette aspiration techniques [33-35], as well as optical and magnetic tweezers measurements [36,37] and experiments that utilize atomic force microscopes [38-45], provide interesting data, but in a contexts that are different from those that we want to simulate);
- cells adapt their shape both to environment and to internal events (the most notable of which is obviously the mitosis)
- cells live in a high-viscosity environment where both the internal viscosity of cells and the external viscosity of the extracellular environment matter.

Clearly, phenomenological approximations and parameterizations carry even more weight in this context. We approximate the mechanical behavior of cells as follows:

- cells are (soft) spheres, and are characterized by the position of their center (a 3D vector) and by their radius;
- we assume that cells are close-packed, and this leads to the introduction of an effective radius that is important in repulsive interactions with other cells;
- adhesion forces are usually very short-ranged, but here we assume a much larger range, and in this way we account for the cell's extensibility and deformation;
- we assume that cells are surrounded by viscous material similar to that found in the extracellular matrix [46];

In the following paragraphs we explain all this in further detail.

## Cell-cell forces

We use a parameterization of the cell-cell forces that is somewhat similar to that used in references [47] and [48]. For small deformations of the cell

membrane, it is reasonable to assume that it behaves as in the Hertz problem (interaction of two spherical membranes) or as in the Boussinesq problem (axisymmetric pressure on a flat membrane) [49,50]: in both cases the force is proportional to  $k_C |x|^{3/2}$ , where x is the relative deviation from the equilibrium position and  $k_c$  is a constant related to the problem parameters. In the Hertz problem the intensity of the force is

$$F_{12} = \frac{\sqrt{R_1 R_2} \left(R_1 + R_2\right)}{\frac{3}{4} \left(\frac{1 - v_1^2}{E_1} + \frac{1 - v_2^2}{E_2}\right)} \left(\frac{R_1 + R_2 - d}{R_1 + R_2}\right)^{3/2}$$
(5.45)

where R denotes the radius, d is the distance of the two spheres' centers, E is Young's modulus, v is Poisson's ratio, and the subscripts denote sphere 1 or 2, so that in this case

$$k_{C} = \frac{\sqrt{R_{1}R_{2}} \left(R_{1} + R_{2}\right)}{\frac{3}{4} \left(\frac{1 - v_{1}^{2}}{E_{1}} + \frac{1 - v_{2}^{2}}{E_{2}}\right)}$$
(5.46)

and

$$x = \frac{d - (R_1 + R_2)}{(R_1 + R_2)}$$
(5.47)

Under compression the repulsive force (S.45) can be quite large, especially after mitosis: such large forces are not observed in real cells, and thus we assume that at small enough separations *d* the force flattens out and has a constant modulus. We set the position of this flattening so that we obtain the observed duration of mitosis.

When cells are drawn apart the force is attractive, because of adhesion molecules on the cell membrane [51]. The force range of the adhesion molecules is quite small (tens of nanometers) but here we assume a far larger range, as

large as a few microns: in this way we account - albeit phenomenologically - for shape deformations in the case of attractive forces. Cell-cell adhesion depends on the number of links between adhesion molecules on both cell membranes [52], and is a probabilistic process: if we assume a roughly Gaussian probability density which depends on the relative deviation *x*, then

$$p(x) = \frac{1}{\sqrt{2\pi\sigma^2}} \exp\left(-\frac{(x-x_0)^2}{2\sigma^2}\right)$$
(5.48)

is the probability density that a link is detached at relative distance x, where  $x_0$ and  $\sigma^2$  are parameters that must be adjusted. This means that the average number of detached links at relative distance x is proportional to the cumulative probability

$$P(x) = \int_{0}^{x} p(x') dx' = \int_{0}^{x} \frac{1}{\sqrt{2\pi\sigma^{2}}} \exp\left(-\frac{(x'-x_{0})^{2}}{2\sigma^{2}}\right) dx' \approx \frac{1}{2} \left[1 + \operatorname{erf}\left(\frac{x-x_{0}}{\sqrt{2}\sigma}\right)\right]$$
(5.49)

Computing the error function in the simulation program leads to an increase of the computational load, and thus we replace the expression (S.49) with the approximate formula

$$P(x) \approx \frac{1}{2} \left[ 1 + \tanh\left(\sqrt{\frac{2}{\pi\sigma^2}} \left(x - x_0\right)\right) \right]$$
(5.50)

(the hyperbolic tangent approximates the error function everywhere to better than 2% precision). Finally we take a force that is proportional both to the number of links and to the previously calculated shape  $k_C |x|^{3/2}$ :

$$F_{a}(x) \approx \frac{1}{2} \left[ 1 + \tanh\left(\sqrt{\frac{2}{\pi\sigma^{2}}} (x - x_{0})\right) \right] k_{c} |x|^{3/2}$$
(5.51)

The values of the parameters E, v,  $\sqrt{\frac{2}{\pi\sigma^2}}$ , and  $x_0$  are listed in table S4 and the modulus of the force is shown in figure S4; the direction of the force is always along the line that connects the cell centers.

#### Brownian motion

Cells are small and Brownian motion, i.e. random molecular forces, could play an important role, however in the high viscosity medium of a cell spheroid we find that Brownian motion is negligible. It is easy to estimate this from the 1D Langevin equation for Brownian motion along *x*:

$$m\frac{d^2x}{dt^2} = -\gamma\frac{dx}{dt} + n(t)$$
(5.52)

where *m* is the cell's mass,  $\gamma$  is the friction coefficient, and n(t) is a white noise process which is associated to the molecular thermal motion. It can easily be shown that the mean square fluctuation [53] is

$$\left\langle x^{2}\left(t\right)\right\rangle = 2\frac{kT}{\gamma}t\tag{5.53}$$

and does not depend on the particle mass. Now if we take cells in a highviscosity environment,  $\eta \approx 1 \text{ Pa} \cdot \text{s}$ , and estimate the friction coefficient using the Stokes-Einstein formula for a cell with radius  $r \approx 5 \ \mu\text{m}$ , we find  $\gamma = 6\pi \eta r \approx 9.4 \cdot 10^{-5} \text{ N} \cdot \text{s/m}$ , then at  $T \approx 300 \text{ K}$  the mean square fluctuation and the RMS fluctuation are

$$\langle x^{2}(t) \rangle \approx (8.8 \cdot 10^{-17} \text{ m}^{2}/\text{s}) t$$
 (5.54a)

$$\sqrt{\langle x^2(t) \rangle} \approx \left(9.4 \ nm/\sqrt{s}\right) \sqrt{t}$$
 (5.54b)

so that it takes a time  $t \approx 10^4$  s for an RMS fluctuation of 1 µm to develop. It follows that the Brownian motion of cells is negligible in the case a simulation of tumor spheroids with reasonable time steps (the time step cannot in any case be longer than the duration of the shortest phase, which is the M-phase and lasts ~ 2000 s. moreover time steps should be shorter than ~ 100 s to provide a good sampling schedule of most biochemical processes).

## Dynamical equations

The dynamical equations for the motion of cells are similar to those of dissipative particle dynamics [54,55], without stochastic terms (since Brownian motion is negligible):

$$m_n \frac{d\mathbf{v}_n}{dt} = -\gamma \mathbf{v}_n - \sum_k \gamma_{n,k} \frac{(\mathbf{v}_n - \mathbf{v}_k) \cdot (\mathbf{r}_n - \mathbf{r}_k)}{|\mathbf{r}_n - \mathbf{r}_k|^2} (\mathbf{r}_n - \mathbf{r}_k) + \mathbf{F}_n + \sum_k \mathbf{F}_{n,k}$$
(S.55)

where the indices *n* and *k* denote cells, the sums  $\sum_{k}$  are over all neighboring cells,  $m_n$  denotes the mass,  $\mathbf{r}_n$  and  $\mathbf{v}_n$  are position and velocity vectors of the cell,  $\gamma$  is the environmental friction coefficient,  $\gamma_{n,k}$  is the friction between the *n*-th and the *k*-th cell,  $\mathbf{F}_n$  is an external force (like gravity), and finally  $\mathbf{F}_{n,k}$  is the force that cell *k* exerts on cell *n*.

## Analysis of the simplified dynamical equations

The equations (S.55) include a cell-cell friction that acts only along the line that joins the cells' centers, and if we relax this constraint, the equations become somewhat simpler and easier to analyze:

$$m_n \frac{d\mathbf{v}_n}{dt} = -\gamma \mathbf{v}_n - \sum_k \gamma_{n,k} \left( \mathbf{v}_n - \mathbf{v}_k \right) + \mathbf{F}_n + \sum_k \mathbf{F}_{n,k}$$
(5.56)

If there are N cells, then the system of equations (S.56) is a linear system of 3N ordinary differential equations, and we can solve it using an implicit-explicit Euler method [56,57] (implicit over velocities, explicit over intercellular forces which are assumed to change slowly), so that the discretized version of the differential equation now writes

$$\mathbf{v}_{n}(t+\Delta t) = \mathbf{v}_{n}(t) + \left\{ -\frac{\gamma}{m_{n}} \mathbf{v}_{n}(t+\Delta t) - \frac{1}{m_{n}} \sum_{k} \gamma_{n,k} \left( \mathbf{v}_{n}(t+\Delta t) - \mathbf{v}_{k}(t+\Delta t) \right) + \frac{1}{m_{n}} \left[ \mathbf{F}_{n}(t) + \sum_{k} \mathbf{F}_{n,k}(t) \right] \right\} \Delta t$$

$$\mathbf{r}_{n}(t+\Delta t) = \mathbf{r}_{n}(t) + \mathbf{v}_{n}(t+\Delta t) \Delta t$$
(S.57b)

Equations (S.57a) are a linear system that could in principle be solved using one of the standard methods [29], however the computational complexity of the standard exact solutions is  $O((3N)^3)$ , and since N can be of the order of several hundreds of thousands or a few millions of cells, the computational burden of the exact solutions is huge. For this reason we resort to an approximate, iterative method (this holds for the solution of the complete differential system (S.55) as well). Formally solving the equations (S.57a) for  $\mathbf{v}_n(t + \Delta t)$ , we find

$$\mathbf{v}_{n}(t+\Delta t) = \frac{\mathbf{v}_{n}(t) + \left\{\sum_{k} \gamma_{n,k} \mathbf{v}_{k}(t+\Delta t) + \left[\mathbf{F}_{n}(t) + \sum_{k} \mathbf{F}_{n,k}(t)\right]\right\} \frac{\Delta t}{m_{n}}}{1 + \left(\gamma + \sum_{k} \gamma_{n,k}\right) \frac{\Delta t}{m_{n}}}$$
(5.58)

and we can rewrite this as

$$\mathbf{v}_{n}^{(i+1)}(t+\Delta t) = \frac{\mathbf{v}_{n}(t) + \left\{\sum_{k} \gamma_{n,k} \mathbf{v}_{k}^{(i)}(t+\Delta t) + \left[\mathbf{F}_{n}(t) + \sum_{k} \mathbf{F}_{n,k}(t)\right]\right\} \frac{\Delta t}{m_{n}}}{1 + \left(\gamma + \sum_{k} \gamma_{n,k}\right) \frac{\Delta t}{m_{n}}}$$

$$= \frac{\sum_{k} \gamma_{n,k} \mathbf{v}_{k}^{(i)}(t+\Delta t) \frac{\Delta t}{m_{n}}}{1 + \left(\gamma + \sum_{k} \gamma_{n,k}\right) \frac{\Delta t}{m_{n}}} + \frac{\mathbf{v}_{n}(t) + \left[\mathbf{F}_{n}(t) + \sum_{k} \mathbf{F}_{n,k}(t)\right] \frac{\Delta t}{m_{n}}}{1 + \left(\gamma + \sum_{k} \gamma_{n,k}\right) \frac{\Delta t}{m_{n}}}$$
(5.59)

where the index *i* identifies the *i*-th iteration. It is easy to see that the iterative algorithm converges: we start by subtracting two successive iterations

$$\mathbf{v}_{n}^{(i+2)}(t+\Delta t) - \mathbf{v}_{n}^{(i+1)}(t+\Delta t) = \frac{\sum_{k} \gamma_{n,k} \left[ \mathbf{v}_{k}^{(i+1)}(t+\Delta t) - \mathbf{v}_{k}^{(i)}(t+\Delta t) \right] \frac{\Delta t}{m_{n}}}{1 + \left( \gamma + \sum_{k} \gamma_{n,k} \right) \frac{\Delta t}{m_{n}}}$$
(S.60)

from which we obtain the series of inequalities

$$\max_{n} \left| \mathbf{v}_{n}^{(i+2)}(t+\Delta t) - \mathbf{v}_{n}^{(i+1)}(t+\Delta t) \right| \leq \max_{n} \frac{\sum_{k} \gamma_{n,k} \left| \mathbf{v}_{k}^{(i+1)}(t+\Delta t) - \mathbf{v}_{k}^{(i)}(t+\Delta t) \right| \frac{\Delta t}{m_{n}}}{1 + \left( \gamma + \sum_{k} \gamma_{n,k} \right) \frac{\Delta t}{m_{n}}} \leq \max_{n} \frac{\sum_{k} \gamma_{n,k} \frac{\Delta t}{m_{n}}}{1 + \left( \gamma + \sum_{k} \gamma_{n,k} \right) \frac{\Delta t}{m_{n}}} \cdot \max_{n} \left| \mathbf{v}_{k}^{(i+1)}(t+\Delta t) - \mathbf{v}_{k}^{(i)}(t+\Delta t) \right| \\ = c \cdot \max_{n} \left| \mathbf{v}_{k}^{(i+1)}(t+\Delta t) - \mathbf{v}_{k}^{(i)}(t+\Delta t) \right| \tag{5.61}$$

and convergence follows when we notice that

$$c = \max_{n} \frac{\sum_{k} \gamma_{n,k} \frac{\Delta t}{m_{n}}}{1 + \left(\gamma + \sum_{k} \gamma_{n,k}\right) \frac{\Delta t}{m_{n}}} < 1$$
(S.62)

i.e., the maximum difference  $\max_{n} |\mathbf{v}_{k}^{(i+1)}(t + \Delta t) - \mathbf{v}_{k}^{(i)}(t + \Delta t)|$  between successive iterations approaches zero for large *i*.

In the section on Brownian motion we have already taken a rather high environmental viscosity  $\eta \approx 1 \text{ Pars}$ , which corresponds to the viscosity of the hyaluronate and collagen in extracellular spaces [58,59]; the viscosity of cell-cell interactions is even higher [31-45], of the order of  $\eta \approx 200 \text{ Pars}$ . If we take cells with radius  $r \approx 5 \text{ µm}$ , so that  $m \approx 5 \cdot 10^{-13} \text{ kg}$ , we note that  $\gamma \approx 9.4 \cdot 10^{-5} \text{ N} \cdot \text{s/m}$  and that  $\sum_{k} \gamma_{n,k} \approx 0.02 \text{ N} \cdot \text{s/m}$ , and we let  $\Delta t = 100 \text{ s}$ , then we find  $c \approx 0.995$ . Using this value we can estimate the worst-case convergence speed: from equation (S.62) we see that we achieve a ten-fold accuracy improvement after *s* steps, so that  $c^s = 1/10$ , i.e.,  $s = -1/\log_{10} c \approx 500$ . We remark here that this estimate depends critically on the environmental viscosity: we assume the viscosity of hyaluronate, but if we took a lower viscosity environment, like water with  $\eta \approx 0.001 \text{ Pars}$ , we would find  $\gamma \approx 9.4 \cdot 10^{-8} \text{ N} \cdot \text{s/m}$  and  $c \approx 0.999995$ , so that  $s \approx 500000$ . In the case of a water environment, convergence of the iterative algorithm is so slow that it becomes totally impractical, but fortunately we can stick to the other, more favourable case.

## Solution of the complete nonlinear system of dynamical equations

The simplified equations (S.56) are useful to analyze algorithmic convergence, but the program utilizes the full equations (S.55). In this case the equations obtained from the implicit-explicit Euler method for the *x*-component of the velocity vector are

$$\begin{aligned} v_{n,x}(t+\Delta t) \Biggl[ 1 + \frac{\Delta t}{m_n} \Biggl( \gamma + \sum_k \gamma_{n,k} \frac{(r_{n,x} - r_{k,x})}{|\mathbf{r}_n - \mathbf{r}_k|^2} (r_{n,x} - r_{k,x}) \Biggr) \Biggr] \\ + v_{n,y}(t+\Delta t) \frac{\Delta t}{m_n} \sum_k \gamma_{n,k} \frac{(r_{n,y} - r_{k,y})}{|\mathbf{r}_n - \mathbf{r}_k|^2} (r_{n,x} - r_{k,x}) \\ + v_{n,z}(t+\Delta t) \frac{\Delta t}{m_n} \sum_k \gamma_{n,k} \frac{(r_{n,z} - r_{k,z})}{|\mathbf{r}_n - \mathbf{r}_k|^2} (r_{n,x} - r_{k,x}) \\ = v_{n,x}(t) + \frac{\Delta t}{m_n} \Biggl[ \sum_k \gamma_{n,k} \frac{\mathbf{v}_k (t+\Delta t) (\mathbf{r}_n - \mathbf{r}_k)}{|\mathbf{r}_n - \mathbf{r}_k|^2} (r_{n,x} - r_{k,x}) + F_{n,x} + \sum_k F_{n,k;x} \Biggr] \end{aligned}$$
(5.63)

(similar equations hold for the y and z component of the velocity). These equations can also be put into matrix form

$$\mathbf{M} \cdot \mathbf{v}_n \left( t + \Delta t \right) = \mathbf{b} \tag{5.64}$$

where

$$\mathbf{b} = \mathbf{v}_{n}(t) + \frac{\Delta t}{m_{n}} \left( \sum_{k} \gamma_{n,k} \frac{\mathbf{v}_{k}(t + \Delta t) \cdot (\mathbf{r}_{n} - \mathbf{r}_{k})}{\left|\mathbf{r}_{n} - \mathbf{r}_{k}\right|^{2}} (\mathbf{r}_{n} - \mathbf{r}_{k}) + \mathbf{F}_{n} + \sum_{k} \mathbf{F}_{n,k} \right)$$
(S.65)

and

$$\mathbf{M} = \begin{pmatrix} 1 + \frac{\Delta t}{m_n} \left( \gamma + \sum_k \gamma_{n,k} \frac{\left(r_{n,x} - r_{k,x}\right)^2}{\left|\mathbf{r}_n - \mathbf{r}_k\right|^2} \right) & \frac{\Delta t}{m_n} \sum_k \gamma_{n,k} \frac{\left(r_{n,y} - r_{k,y}\right)}{\left|\mathbf{r}_n - \mathbf{r}_k\right|^2} \left(r_{n,x} - r_{k,x}\right) & \frac{\Delta t}{m_n} \sum_k \gamma_{n,k} \frac{\Delta t}{m_n - \mathbf{r}_k} \left(r_{n,x} - r_{k,x}\right) + \frac{\Delta t}{m_n - \mathbf{r}_k} \left(r_{n,y} - r_{k,y}\right)^2 \\ \frac{\Delta t}{m_n - \mathbf{r}_k} \sum_k \gamma_{n,k} \frac{\left(r_{n,x} - r_{k,x}\right)}{\left|\mathbf{r}_n - \mathbf{r}_k\right|^2} \left(r_{n,y} - r_{k,y}\right) & 1 + \frac{\Delta t}{m_n} \left(\gamma + \sum_k \gamma_{n,k} \frac{\left(r_{n,y} - r_{k,y}\right)^2}{\left|\mathbf{r}_n - \mathbf{r}_k\right|^2}\right) & \frac{\Delta t}{m_n - \mathbf{r}_k} \sum_k \gamma_{n,k} \frac{\left(r_{n,z} - r_{k,z}\right)}{\left|\mathbf{r}_n - \mathbf{r}_k\right|^2} \left(r_{n,z} - r_{k,z}\right) \\ \frac{\Delta t}{m_n - \mathbf{r}_k} \sum_k \gamma_{n,k} \frac{\left(r_{n,x} - r_{k,x}\right)}{\left|\mathbf{r}_n - \mathbf{r}_k\right|^2} \left(r_{n,z} - r_{k,z}\right) & \frac{\Delta t}{m_n - \mathbf{r}_k} \sum_k \gamma_{n,k} \frac{\left(r_{n,z} - r_{k,z}\right)}{\left|\mathbf{r}_n - \mathbf{r}_k\right|^2} \right) \\ \frac{\Delta t}{m_n - \mathbf{r}_k} \sum_k \gamma_{n,k} \frac{\left(r_{n,z} - r_{k,z}\right)}{\left|\mathbf{r}_n - \mathbf{r}_k\right|^2} \left(r_{n,z} - r_{k,z}\right) & \frac{\Delta t}{m_n - \mathbf{r}_k} \sum_k \gamma_{n,k} \frac{\left(r_{n,z} - r_{k,z}\right)}{\left|\mathbf{r}_n - \mathbf{r}_k\right|^2} \right) \\ \frac{\Delta t}{m_n - \mathbf{r}_k} \sum_k \gamma_{n,k} \frac{\left(r_{n,z} - r_{k,z}\right)}{\left|\mathbf{r}_n - \mathbf{r}_k\right|^2} \left(r_{n,z} - r_{k,z}\right) & \frac{\Delta t}{m_n - \mathbf{r}_k} \sum_k \gamma_{n,k} \frac{\left(r_{n,z} - r_{k,z}\right)}{\left|\mathbf{r}_n - \mathbf{r}_k\right|^2} \right) \\ \frac{\Delta t}{m_n - \mathbf{r}_k} \sum_k \gamma_{n,k} \frac{\left(r_{n,z} - r_{k,z}\right)}{\left|\mathbf{r}_n - \mathbf{r}_k\right|^2} \left(r_{n,z} - r_{k,z}\right) & \frac{\Delta t}{m_n - \mathbf{r}_k} \sum_k \gamma_{n,k} \frac{\left(r_{n,z} - r_{k,z}\right)}{\left|\mathbf{r}_n - \mathbf{r}_k\right|^2} \right) \\ \frac{\Delta t}{m_n - \mathbf{r}_k} \sum_k \gamma_{n,k} \frac{\left(r_{n,z} - r_{k,z}\right)}{\left|\mathbf{r}_n - \mathbf{r}_k\right|^2} \left(r_{n,z} - r_{k,z}\right) + \frac{\Delta t}{m_n - \mathbf{r}_k} \left(r_{n,z} - r_{k,z}\right)} \left(r_{n,z} - r_{k,z}\right) \right)$$

The formal solution of the equations (S.64) is

$$\mathbf{v}_{n}(t + \Delta t) = \mathbf{M}^{-1} \cdot \mathbf{b}$$

$$= \mathbf{M}^{-1} \cdot \left[ \mathbf{v}_{n}(t) + \frac{\Delta t}{m_{n}} \left( \sum_{k} \gamma_{n,k} \frac{\mathbf{v}_{k}(t + \Delta t) \cdot (\mathbf{r}_{n} - \mathbf{r}_{k})}{|\mathbf{r}_{n} - \mathbf{r}_{k}|^{2}} (\mathbf{r}_{n} - \mathbf{r}_{k}) + \mathbf{F}_{n} + \sum_{k} \mathbf{F}_{n,k} \right) \right]$$
(S.67)

and this equation can be recast in the iterative form

$$\mathbf{v}_{n}^{(i+1)}(t+\Delta t) = \mathbf{M}^{-1} \cdot \mathbf{b}$$

$$= \mathbf{M}^{-1} \cdot \left[ \mathbf{v}_{n}(t) + \frac{\Delta t}{m_{n}} \left( \sum_{k} \gamma_{n,k} \frac{\mathbf{v}_{k}^{(i)}(t+\Delta t) \cdot (\mathbf{r}_{n} - \mathbf{r}_{k})}{\left| \mathbf{r}_{n} - \mathbf{r}_{k} \right|^{2}} (\mathbf{r}_{n} - \mathbf{r}_{k}) + \mathbf{F}_{n} + \sum_{k} \mathbf{F}_{n,k} \right) \right]$$
(5.68)

where once again the index *i* identifies the *i*-th iteration.

#### Mitosis

When a cell completes the M-phase it is replaced by two daughter cells. The program selects a random direction for the axis that joins the centers of the daughter cells, and then it computes the positions of their centers. Since the total volume of the daughter cells is equal to the volume of the mother, the radius of each daughter cell is roughly equal to 80% of the radius  $R_0$  of the mother, i.e., the distance between the new centers is about  $0.4R_0$  (see figure S5). Here the centers are only representative of an "average" cell position and are used to compute forces; the new cells are actually compressed and deformed to fit in the original volume.

The distance  $0.4 R_0$  also sets the maximum value of the repulsive force: indeed when the daughter cells are not surrounded by other cells, the distance of their centers must increase from about  $0.4 R_0$  to  $1.6 R_0$  in a time equal to the duration of the M-phase. For cells with radius  $R_0 \approx 5 \ \mu m$  this means that the distance traveled by each cell is about  $3 \ \mu m$ . If we take the total duration of the M-phase about 2000 s, then the average speed of each cell is about 1.5 nm/s, and the

average repulsive force is  $F = \gamma v \approx 30$  pN, taking a cell-cell viscosity  $\eta \approx 200$  Pa·s as above.

# D. List of biophysical constants, parameter values

Γ

basic informa	tions		
name	cell name (identifier)		
mark	optional label to select a subset of cells (i.e., cells with a marker of some kind)		
type	cell phenotype (cell-specific list of parameters)		
Ť	local cell temperature (not used now)		
cell state			
phase	cell phase		
death_condition	label that records the reason of death (valid for dead cells only)		
age	cell age (since birth)		
phase_age	age of present cell phase		
age_mother	age of mother at mitosis		
n_mitosis	number of cell generations since start of simulation		
geometric and	topological informations		
x,y,z	cell position vector		
vx, vy, vz	cell velocity vector		
r	cell radius		
surface	cell surface		
volume	cell volume		
mass	cell mass		
volume_extra	volume of extracellular space surrounding cell		
neign	number of neighbors		
vneign	list of neignboring cells		
valst	distances to neighboring cells		
vcsurj	contact surfaces with neighboring cells		
glik	vector of geometric factors		
contact_surj			
ison(H	boolean variable that is true if the cell is in contact with the environment		
env surf	cell-environment contact area		
env_surj			
metabolic var	iables		
nHi	internal nH		
M	number of mitochondria		
G	mass of alucose inside cell		
G6P	mass of G6P inside cell		
02	mass of oxygen inside cell		
store	mass of glycogen inside cell		
A	mass of glutamine inside cell		
AcL	mass of lactic acid inside cell		
protein	total mass of proteins inside cell		
prot rate	total protein production rate		
DNA	total amount of DNA produced by replication		
DNA_rate	DNA replication rate		
•	· · · ·		

ratos			
Tules	ducase observation rate		
GADSRULE	glucose absorption rate		
AlbeRate	glucose consumption rate		
AADSRULE	glutamine absorption rate		
ACONSKALE	glutamine consumption rate		
StoreFillRale			
	glycogen consumption rate		
AcLRate	lactic acid production rate		
AcLOUTRate	lactic acid expulsion rate		
ATP-related va	ariables		
ATP_St	standard ATP		
ATP_Ox	ATP production rate (by oxidative phosphorylation)		
ATP_NOx	ATP production rate (by anaerobic glycolysis)		
ATP2	ATP production rate from the glycogen store		
ATP3	ATP production rate from glutamine		
ConsATP	total ATP consumption rate due to metabolic activity		
ConsATP 1	ATP consumption rate associated to mitochondrial activity and other volume-		
	dependent processes		
ConsATP 2	ATP consumption rate associated to the production of proteins		
ConsATP 3	ATP consumption rate associated to DNA production		
ConsATP 4	(reserved for future use)		
ConsATP 5	ATP consumption rate associated to mitochondrial proliferation		
ATPtot	total ATP variation rate (this is the sum of all production and consumption		
	rates)		
ATPp	ATP pool (total mass of ATP inside cell)		
ATPmin	minimum tolerable ATPp level (it depends on the size of cell: larger cells have a larger ATPmin level)		
ATPstart	this variable stores the initial value of ATPp in the newly born cell		
ATPprod	ATP produced in cell during the last simulation step		
ATPcons	ATP consumed in cell during the last simulation step		
variables that	characterize the extracellular space around a cell		
pН	pH in extracellular space		
G_extra	mass of glucose in extracellular space		
	mass of glutamine in extracellular space		
AcL_extra	mass of lactic acid in extracellular space		
other variables			
SensO?	fraction of oxygen available with respect to request		
ConsO	oxygen consumption rate		
protoing and DNA			
DNA sproad	individual variation of DNA production rate (this is a quite rough modeling of		
DNA_spreaa	variable DNA synthesis because of individual DNA damages)		
M_T	average duration of M-phase		
pRb	total pRb mass inside cell		
ConcS	molar concentration of substrate needed for thresholds		
cyclinD	total cyclin D mass inside cell		
cyclinE	total cyclin E mass inside cell		
cyclinX	total mass of A and B cyclins inside cell		
NpRbk	total number of active pRb molecules		

Table S1: list of cell status variables in the simulation program. (ctd.)

Table S2 : list of the biophysical constants that are hard-coded in the program

Name of constant Value		Description
N <sub>AV</sub>	6.022e23	Avogadro constant
Faraday	96485.34 Coulomb	Faraday constant
PM <sub>G</sub>	0.18 kg	Molecular weight of glucose
PM <sub>ATP</sub>	0.507 kg	Molecular weight of ATP
PM <sub>02</sub>	0.032 kg	Molecular weight of $O_2$
PM <sub>AcL</sub>	0.090 kg	Molecular weight of lactate
PM <sub>G6P</sub>	0.270 kg	Molecular weight of G6P
PM <sub>pRb</sub>	110. kg	Molecular weight of pRb
PMA	0.146 kg	Molecular weight of glutamine
PM <sub>prot</sub>	66.476 kg	Average molecular weight of proteins
PM <sub>cyclinD</sub>	33.729 kg	Molecular weight of cycline D
PM <sub>cyclinE</sub>	47.077 kg	Molecular weight of cycline E
PM <sub>cyclinX</sub>	52. kg	Molecular weight of cycline A+B
Gibbs <sub>ATP</sub>	3.1.10 <sup>4</sup> J/mol	Gibbs free energy from ATP hydrolysis

The following tables list the parameters that represent the the environment and the cell phenotype, and that the program reads from external files, and the diffusion and viscosity constants, that are still hard-coded in the simulation program:

Table S3: Starting values for the environmental concentrations (standard
atmospheric pressure and culture medium)

Substance	Environmental
	concentration
<i>O</i> <sub>2</sub>	0.007 kg/m³
glucose	0.9 kg/m³
glutamine	0.4 kg/m <sup>3</sup>

Parameter name	Value	Notes and references
VMAX1	$2 \cdot 10^{9} \text{ kg} \cdot \text{s}^{-1} \cdot \text{m}^{-2}$	[60].
VMAX2	1.2·10 <sup>-19</sup> kg·s <sup>-1</sup>	[61].
VMAX22	$1.2 \cdot 10^{-18} \text{ kg} \cdot \text{s}^{-1}$	[61].
VMAXA	$1 \cdot 10^{-9} \text{ kg} \cdot \text{s}^{-1} \cdot \text{m}^{-2}$	[62].
VMAXP	3.10 <sup>-19</sup> Kg·s <sup>-1</sup>	Estimated in this work by considering that the average protein
	5	content of a cell is approx. $2.6 \cdot 10^{14} \text{ kg}$
VMAXP <sub>4</sub>	1.018·10 <sup>-20</sup> kg·s <sup>-1</sup>	Estimated in this work by considering that glutamine mass is
	5	approx. 5% of the total protein mass in the cell. See also [63].
VMAXPATP	1.81·10 <sup>-18</sup> kg·s <sup>-1</sup>	Estimated in this work by considering albumin (585
	5	aminoacids, 584 aminoacid bonds) as the "average protein" and
		that 2 molecules of ATP are required for each bond.
VMAXDNA	5·10⁻⁵ molecules/s	Estimated in this work taking 3.10° base pairs/cell and
		assuming that on average the S-phase lasts 25000 s
VMAXDNA₄	5.847·10 <sup>-20</sup> kg·s <sup>-1</sup>	Estimated in this work by considering that the entire
	5	duplication of the DNA requires $6.10^{\circ}$ glutamine molecules. See
		also [64].
VMAXDNA ΑΤΡ	9.01.10 <sup>-20</sup> kg·s <sup>-1</sup>	Estimated in this work by considering that, on average, the free
All		energy of each phosphodiester bond in DNA is 5.3 kcal/mole.
		See. e.g., [65].
VMAXM	$2 \cdot 10^{-3}$	Estimated in this work, assuming an average number of 150
	mitochondria/s	mitochondria per cell. See. e.g., [66.67].
VMAXM	$7.3 \cdot 10^{-23} \text{ kg} \cdot \text{s}^{-1}$	Estimated in this work by considering that, on average, there
		are 1000 molecules of mitochondrial DNA/cell and that each
		molecule is 16000 bp. See. e.g., [66].
<i>VMAXM</i> ΑΤΡ	1, 125 · 10 <sup>-22</sup> kg · s <sup>-1</sup>	Estimated in this work as for $VMAXDNA_{4TP}$ . See above.
K1	0.27024 kg⋅m <sup>-3</sup>	[60].
K <sub>2</sub>	1.80 kg $\cdot m^{-3}$	[61].
K22	1.80·10 <sup>-2</sup> kg·m <sup>-3</sup>	[61].
KmA	$0.023798 \text{ kg} \cdot \text{m}^{-3}$	[68].
Ка	0.054 kg · m <sup>-3</sup>	Derived from fitting the model to data in [69]. See also [1].
Ктс	0.096 kg·m <sup>-3</sup>	Derived from fitting the model to data in [69]. See also [1].
Kmd	1.8·10 <sup>-2°</sup> kg·m <sup>-3</sup>	Derived from fitting the model to data in [69]. See also [1].
KmO2	$7.10^{-4}$ kg·m <sup>-3</sup>	Derived from fitting the model to data in [69]. See also [1].
Ктр	6.7·10 <sup>-3°</sup> kg·m <sup>-3</sup>	See ref. 2. In the present case, however, the value is expressed
,	U	in concentration units.
KmDNA	$4.5 \cdot 10^4 (kg \cdot m^{-3})^2$	See [2]. In the present case, however, the value is expressed in
		concentration units.
KmM	$1.46 \cdot 10^{-2} (kg \cdot m^{-3})^2$	Estimated in this work as for $VMAXM_A$ (see above). This value
		is calculated at equilibrium for standard concentrations.
coeffg1	$7.5 \cdot 10^{-3} \text{ s}^{-1}$	Derived from fitting the model to data in [69]. See also [1].
coeffg2	1.08·10 <sup>-3</sup> s <sup>-1</sup>	Derived from fitting the model to data in [69]. See also [1].
coeffg3	6.7·10 <sup>-4</sup> s-1	Derived from fitting the model to data in [69]. See also [1].
coeffr1	3·10 <sup>-20</sup> kg·s-1	Derived from fitting the model to data in [69]. See also [1].
ATP <sub>st</sub>	2.3·10 <sup>-18</sup> kg·s-1	Derived from fitting the model to data in [69]. See also [1].
$V_{min}$	$0.9 \cdot 10^{-16} \text{ m}^{-3}$	Derived from data in [70] considering that the cell nucleus
		occupies between 4% and 14% of cell volume. See also [1].
DVap	$3 \cdot 10^{-6} \text{ s}^{-1}$	Estimated in this work. See also [11].
VMAXACL	9.58·10 <sup>°8</sup> kg·s <sup>-1</sup> ·m <sup>-2</sup>	[71-73].
KmAcL	0.40536 kg·m³	[71-73].
M_T_MEAN	1800. s	[1].
Vwork	1.5·10 <sup>2</sup> kg· s <sup>-1</sup> ·m <sup>-3</sup>	Estimated in this work by fitting of tumor cell growth data

Table S4: Phenotype (cell parameters)

Table S4: Phenotype (cell parameters). (ctd.)

Parameter name	Value	Notes and references

k pRb	10	[17].
N nDb	16	
Ν_ΡΚΟ	10	[17].
pRb ONOFFratio	1·10 <sup>°</sup>	[74].
nPh fraction	$1.5 \cdot 10^{-2}$	This paper by data fitting See also [2]
ρκυ_μαςτισπ	1.5.10	This paper by data fitting. See also [2].
cyclinD_fraction	1.10-	This paper by data fitting. See also [2].
cyclinF fraction	$3.10^{-3}$	This paper by data fitting See also [2]
		This paper by data fitting. See also [2].
cyclinx_fraction	1.4·10 <sup>-</sup>	This paper by data fitting. See also [2].
Concs 0	$1 \cdot 10^{-3} \text{ mol/l}$	[19,20].
Threeh C start	0.9	
11110511_5_51011	0.8	[2].
Thresh S stop	0.05	[2].
k3MM ,-	$1.10^4 \text{ s}^{-1}$	
KJ/M/M		[17,20].
КтММ	1·10° mol/l	19,20 .
NUCLEAR OB L	46	Estimated in this work by fitting of tumor cell growth data
	45	
ClusteringFactor	15	Estimated in this work by fitting of tumor cell growth data
CvcXThr	$0.8 \cdot 10^{-16}  kg$	[2].
$C_1$	$0.074 \text{ m}^3 \text{ km}^{-1}$	Inverse of standard ATD concentration (see above)
<i>C1</i>	0.074 III Kg	niverse of standard ATP concentration (see above)
C2	2·10 '' m '	Average volume of mitochondria. See [75].
0.	$1  2.10^{-5}  \mathrm{m}^{3} \cdot \mathrm{kg}^{-1} \cdot \mathrm{s}^{-1}$	Estimated in this work by fitting of tymor cell growth data
	1.2 10 III kg 3	Estimated in this work by fitting of tanior cett growth data
YoungMod	1000. Pa	Young's modulus. The elastic properties of cells are quite
		variable: according to [38] Young's modulus of living cells
		spans a very large range (1-100 kPa). Here we choose a value
		which is at lowest extreme of this range, and agrees with
		magning and the in raf [11 12 14 15]
		measurements in ref. [41,42,44,45].
PoissonRatio	0.5	Poisson's ratio. This number relates axial and transverse
		strain and lies in the range $(-1, 0, 5)$ [76]. Here we choose the
		struin und ties in the runge (-1,0.5) [70]. Here we choose the
		highest possible value, 0.5, which corresponds to
		incompressible cells (as in [44])
	40 <del>7</del> 0 l - <sup>3</sup>	
aensity	1070. kg·m	[//,/8].
viscositv	200. Pa·s	Cytoplasmatic viscosity, see [34,79].
adhasian ranga	0 E (coll radius)	This work by fitting of types call growth data
aanesion_range	-0.5·(cell raaius)	This work, by fitting of tumor cell growth data.
adhesion decav	2.	This work, by fitting of tumor cell growth data.
packing factor	0 0047	Adimonsional correction that takes into account the
ραεκτης_μετοι	0.9047	Admiensional correction that takes into account the
		interstitial volume (linear packing factor for an arrangement
		of spheres [80, 81])
extension_coeff	1.1	Correction factor that makes cells a little softer (shall be
		removed in a future version of the program with a better
		definition of the Deisson's ratio. This work by fitting of
		definition of the Poisson's ratio). This work by fitting of
		tumor cell growth data.
artvoluma thickness	$1.0.10^{-7}$ m	Thickness of avtracollular space surrounding each call Direct
extvolume_timexness	1.0.10 111	Thickness of extracellular space surrounding each cell. Direct
		measurements from micrographs of tumor spheroids.
artvolume fraction	03	This is the fraction of external volume that is actually
extvolume_fruction	0.5	
		accessible to diffusing molecules [82-84].
thh slope	2 73	This work by fitting of tymor cell growth data See also [2
ipn_stope	2.75	
		13-16].
tnh thr	6 55	This work by fitting of tumor cell growth data See also [2
·p···_··	0.00	12 1/1
		13-16].
tp11 slope	10.91	This work, by fitting of tumor cell growth data. See also [2.
		12 141
		13-10].
tp11 thr	6.9625	This work, by fitting of tumor cell growth data. See also [2,
<i>F</i> =		12,161
- <i>i</i>	- /-	
a2c_slope	2.42	This work, by fitting of tumor cell growth data. See also [2,
- •		13-161
2 (1	( 02	
azc_tnr	6.92	This work, by fitting of tumor cell growth data. See also [2,
		13-161.
c <sup>2</sup> a clana	2 12	This work by fitting of tymer call growth data Sec also [2
czu_siope	2.42	This work, by fitting of tunior cell growth data. see also [2,
		13-16].
cla thr	6 97	This work by fitting of tumor cell growth data See also 12
czu_1111	0.72	1115 TOTK, by fitting of tunior tell growth dutu. See also [2,
		13-10J.
a2cA slope	2.42	This work, by fitting of tumor cell growth data. See also [2]
	· · <b>_</b>	12-161
		וסי-נסן.
a2cA_thr	6.92	This work, by fitting of tumor cell growth data. See also [2.
		13-16]
	0	
clas slong	1.47	This work, by fitting of tumor cell growth data. See also [2.

		13-16].
c2aA_thr	6.92	This work, by fitting of tumor cell growth data. See also [2, 13-16].
Ja2cAcL_slope	1.5	This work, by fitting of tumor cell growth data. See also [2, 13-16].
a2cAcL_thr	7.	This work, by fitting of tumor cell growth data. See also [2, 13-16].
c2aAcL_slope	1.5	This work, by fitting of tumor cell growth data. See also [2, 13-16].
c2aAcL_thr	7.	This work, by fitting of tumor cell growth data. See also [2, 13-16].

Table S5: Other parameters that at present are hard-coded in the program.

Name	Value	Notes and references
Diff_W_G	7.10 <sup>-10</sup> m <sup>2</sup> /s	glucose in water [82].
Diff_ES_G	$2 \cdot 10^{-10} \text{ m}^2/\text{ s}$	glucose in the extracellular spaces [82,85].
Diff_W_A	$3 \cdot 10^{-10} \text{ m}^2/\text{ s}$	glutamine in water [82].
Diff_ES_A	7·10 <sup>11</sup> m²/s	glutamine in the extracellular spaces [86].
Diff_W_AcL	$3 \cdot 10^{-10} \text{ m}^2/\text{ s}$	lactate in water [82].
Diff_ES_AcL	$1.4 \cdot 10^{-11} \text{ m}^2/\text{s}$	lactate in the extracellular spaces [87].
Diff_W_O2	$3.2 \cdot 10^{-9} \text{ m}^2/\text{s}$	oxygen in water [48].
Diff_ES_O2	7.10 <sup>11</sup> m²/s	oxygen in the extracellular spaces. This work, by
		fitting of tumor cell growth data.
VISCOSITY_ENV	1. Pa∙s	viscosity of the extracellular space [58].

# E. Some considerations on the computational complexity of the program

This kind of simulation is a real computational challenge. Although the present version of the program is not completely optimized, it has been designed with computational complexity in mind. On average, the time complexity of each individual time step is  $\approx O(N)$ , where N is the number of cells, and this is determined by the Delaunay triangulation [88-91].

However the number of cells grows exponentially, at least in the initial phase of spheroid growth, and thus the computational complexity of each simulation step

is proportional to  $\exp\left(\frac{t}{T}\ln 2\right)$ , where t is the simulated time and T is the average duplication time of cells. The CPU time  $\Delta t_{CPU}$  required to simulate the time span  $\Delta t_{SM}$  is modeled by the expression

$$\Delta t_{CPU} = \tau_1 \cdot N \ln N + \tau_0 N \tag{5.69}$$

where  $\tau_0$  and  $\tau_1$  are experimental parameters that we determine from actual simulation runs, and where we introduce a slightly superlinear term - suggested by the maximum complexity of Delaunay triangulations in the plane - to account for a superlinear growth term with respect to *N*. Mature spheroids slow down their growth rate and we find that at this later stage  $\Delta t_{CPU}$  is roughly proportional to the number of cells:

$$\Delta t_{CPU} \approx \tau_2 N + \Delta t_0 \tag{S.70}$$

where  $\tau_2$  and  $\Delta t_0$  are yet other parameters that we determine from actual runs. This also means that on the whole the total computing time is roughly proportional to  $N^2$ .

The actual values of  $\tau_0$ ,  $\tau_1$ ,  $\tau_2$ , and  $\Delta t_0$  depend on the simulation parameters: figure S6 shows the CPU time  $t_{CPU}$  (s) it takes for one hour of simulated time vs. *N*, in a run where the precision of the global solution of the problem of diffusion, transport and metabolism is fixed at 1%, and the simulation time step is  $\Delta t_{SIM} = 50$  s. This run has been carried out on a MacPro computer (with an Intel Xeon 3500/W3520 quad processor at 2.66 GHz - and where a single core has been used, the program is not yet multithreaded). Figure S7 shows the total CPU time vs. *N*: fitting this curve we can estimate the total CPU time needed to simulate larger spheroids.

## F. Snapshots from a simulation with the parameters of section E.

The final supplementary figures are higher-resolution pictures of central slices from a simulation with the parameters listed in the previous section. The figures correspond to different simulated times.

- figures S8-10: oxygen concentration (color mapping, blue=low concentration, red=high concentration) and flow (yellow arrows, arrow length proportional to flow intensity). Each figure shows a steady inward flow of oxygen.
- figures S11-14: extracellular glucose concentration (color mapping, blue=low concentration, red=high concentration) and flow (yellow arrows, arrow length proportional to flow intensity). Initially the flow of extracellular glucose points only inward, but glucose diffusion is slow, and as the spheroid grows and dead cells in the core contract, an internal region with an outward flow of extracellular glucose slowly develops.
- figures \$15-17: extracellular glutamine concentration (color mapping, blue=low concentration, red=high concentration) and flow (yellow arrows, arrow length proportional to flow intensity). Glutamine behaves much like glucose, and these snapshots show the glutamine flow before and after the development of the split flow.
- figure S18: lactate concentration (color mapping, blue=low concentration, red=high concentration) and flow (yellow arrows, arrow length proportional to flow intensity). Lactate always flows outward in the simulation: this is a single snapshot taken after both glucose and glutamine have developed their split flow regime.
- figures S19-21: velocity in the plane of the slice (yellow arrows, arrow length proportional to flow intensity). Cells in the core perform complex looping motions, while cells in the viable rim always push outward.

Short movies showing features of the same central slice are also available as additional supplementary information:

- S1: development of the necrotic core (red = live cells, black = dead cells).
- S2: flow of extracellular glucose (same coding as figures S11-14);
- S3: map of projected cell velocities (same coding as figures S19-21).

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